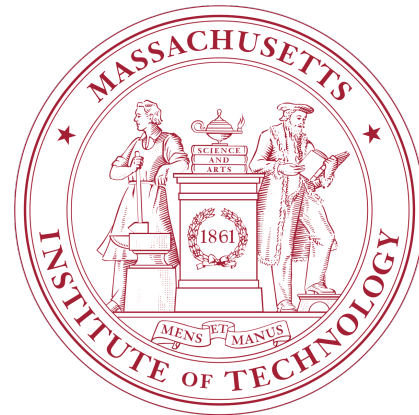


# Loops in the Sky: Recent Applications of Amplitude Methods in Galaxy Clustering and Gravitational Waves

Mikhail (Misha) Ivanov  
CTP / MIT

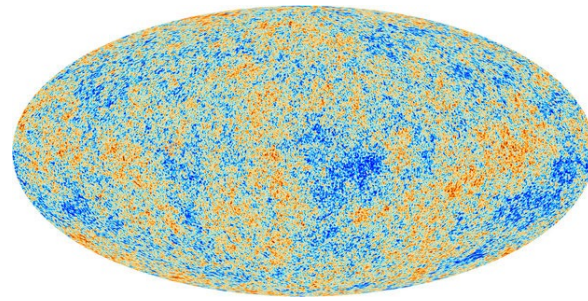
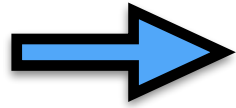


Amplitudes 2024, 10-14 June 2024, IAS

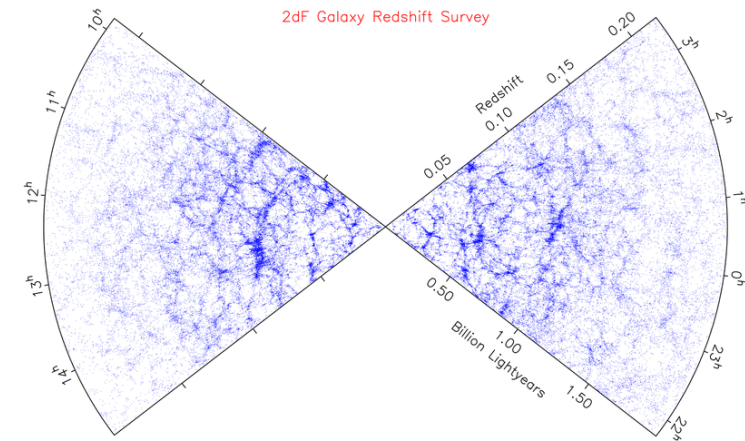
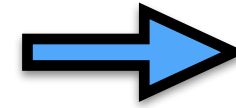
# Cosmology



inflation



CMB



galaxies

$\Lambda$ CDM: Inflation, Cold Dark Matter, Lambda

## Known Unknowns:

What was inflation, exactly?  
Is DM really cold?  
many more ...

Unknown Unknowns:  
Surprises ?

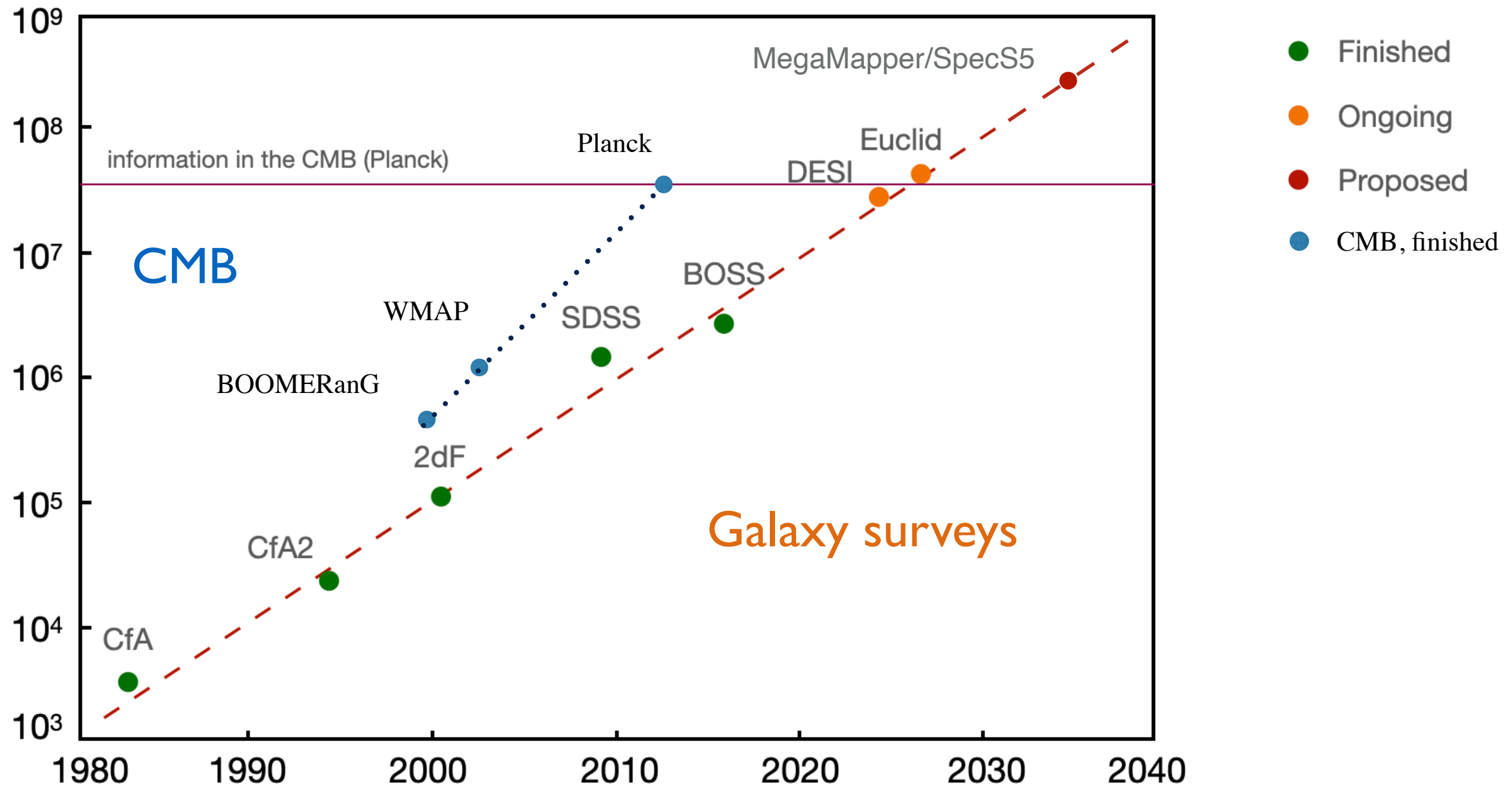


Need to continue measuring fluctuations: galaxies

*also see Hayden Lee's talk*

# Information in Galaxy Surveys

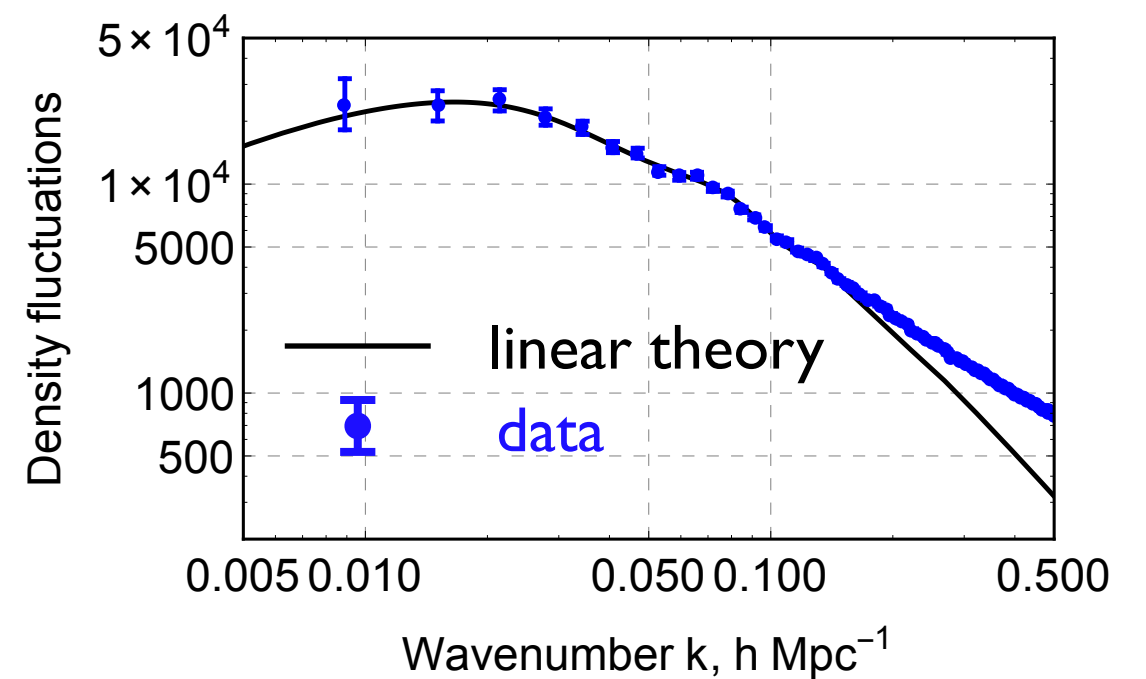
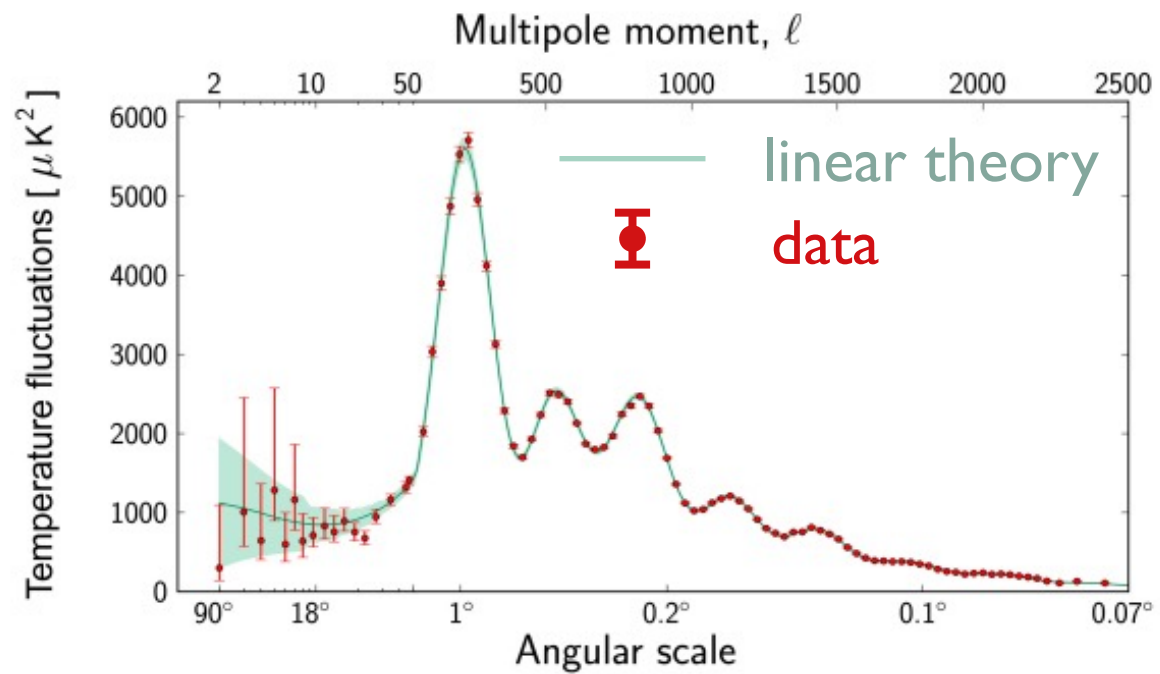
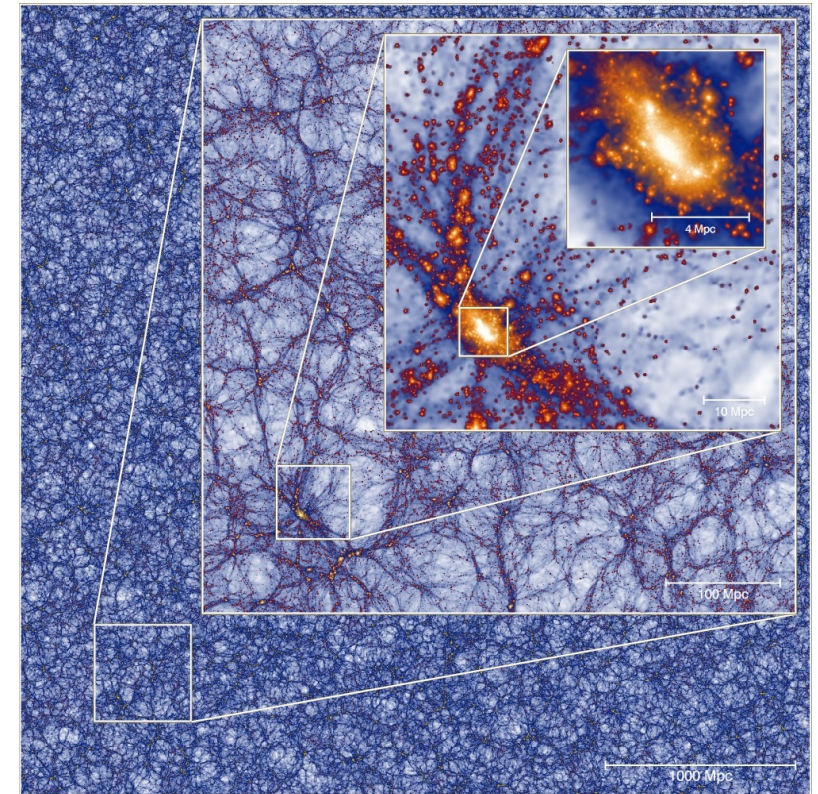
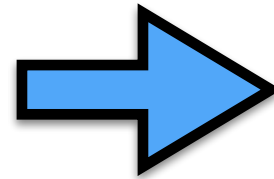
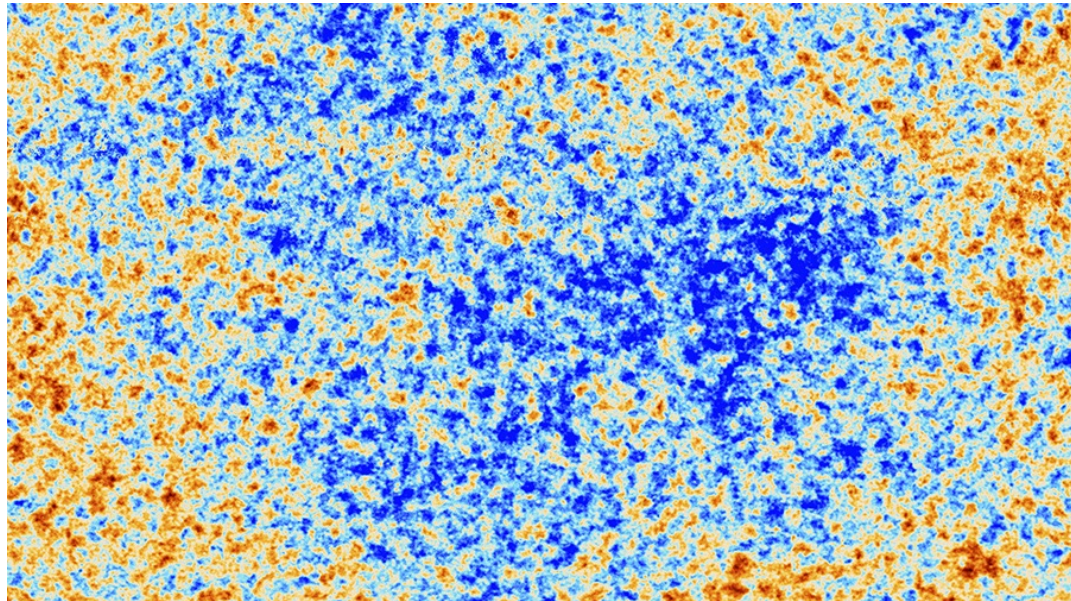
# of modes  $\sim$  information



“No mode left behind” © Nima Arkani-Hamed



# The big problem



non-linearity (non-Gaussianity) is important



# Example: Cosmological Collider in Action



non-Gaussianity probes particle content, interactions, and the speed of propagation



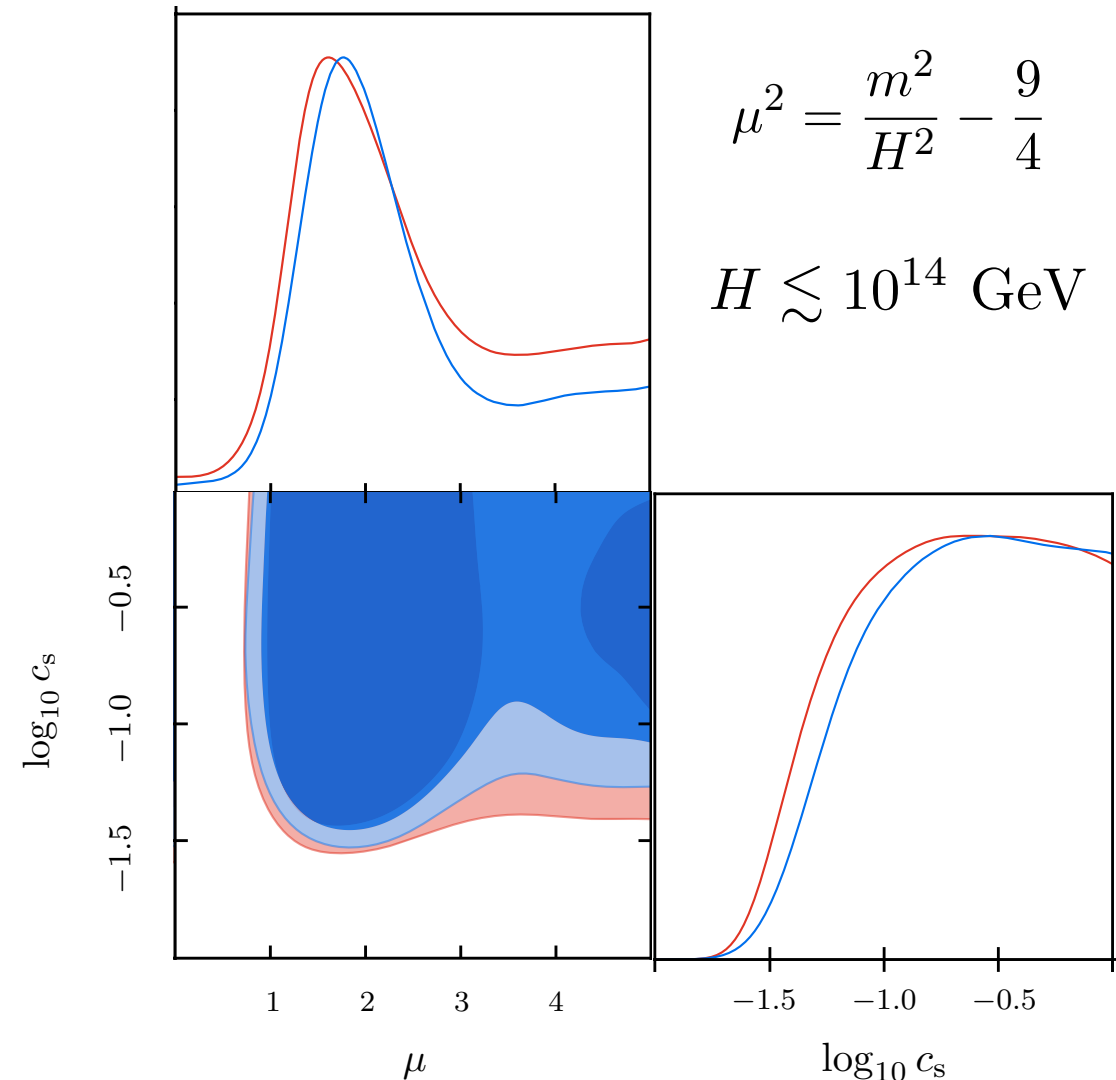
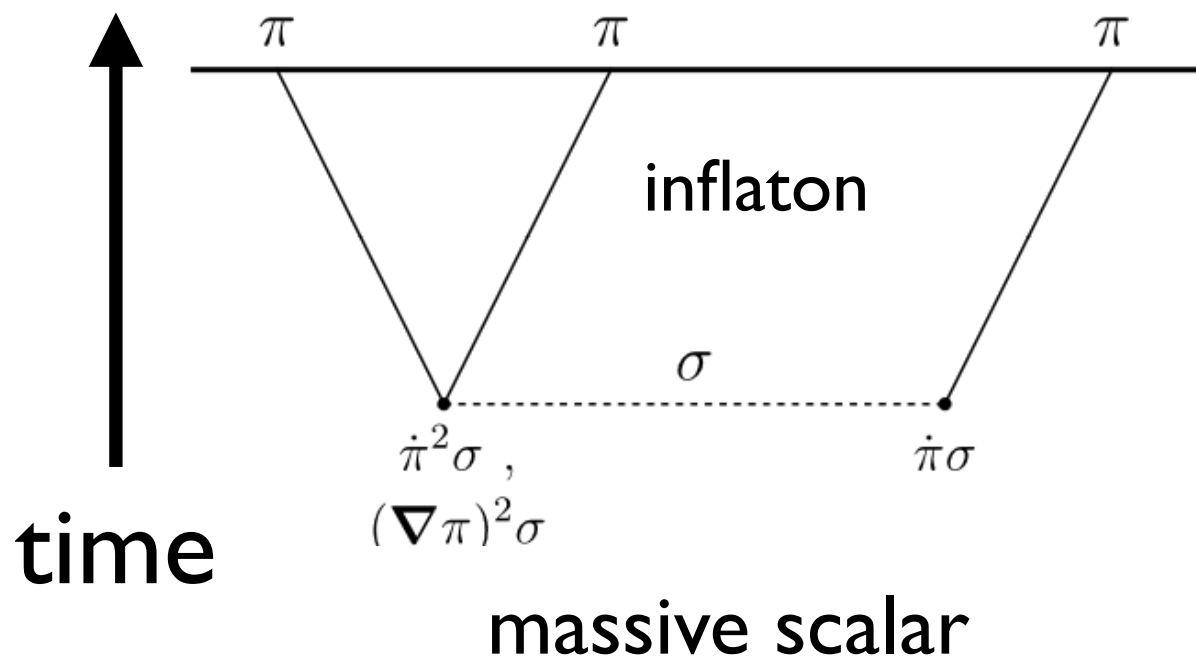
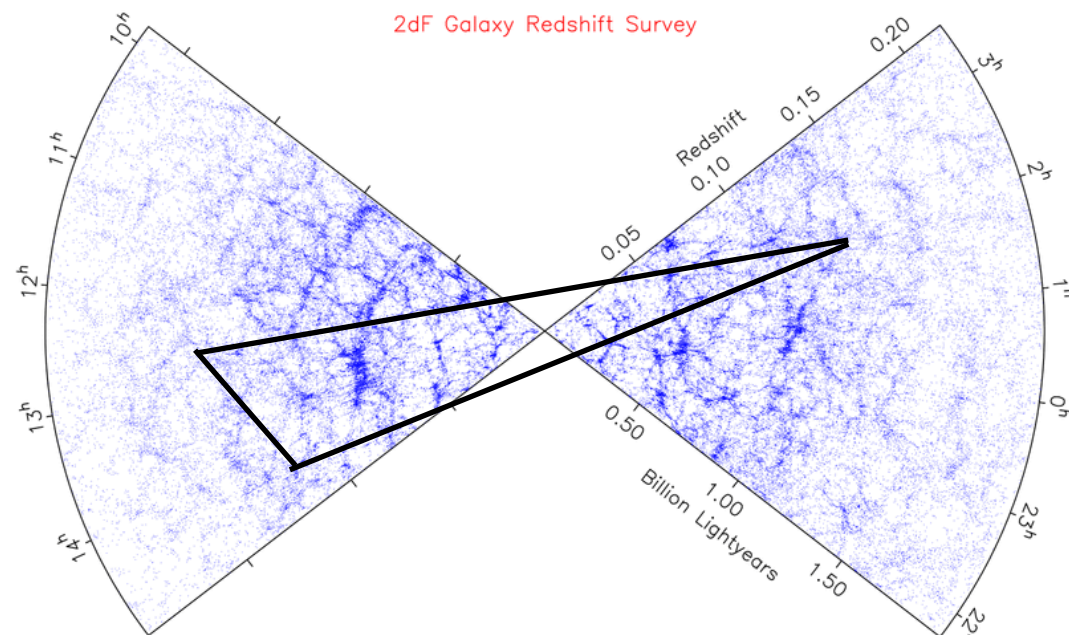
Example: decay of massive particles during inflation



*Cabass, Philcox, MI ++ (2024)*

*Chan, Wang (2009)*

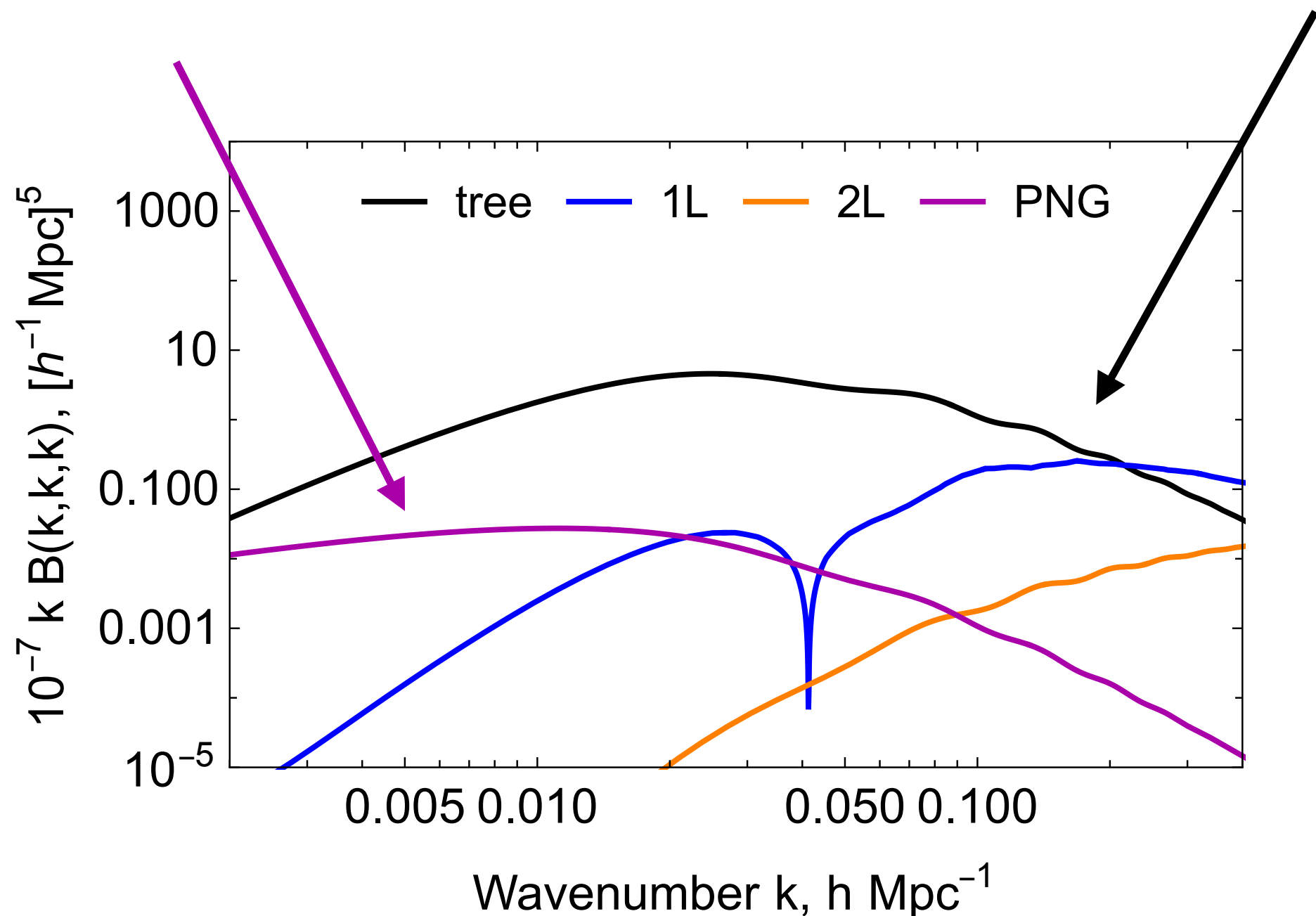
*Arkani-Hamed, Maldacena (2015)*



$$\delta = \frac{\rho}{\langle \rho \rangle} - 1 \quad \langle \delta_g(\mathbf{k}_1) \delta_g(\mathbf{k}_2) \delta_g(\mathbf{k}_3) \rangle = B(k_1, k_2, k_3) \delta_D^{(3)}(\sum \mathbf{k}_i)$$

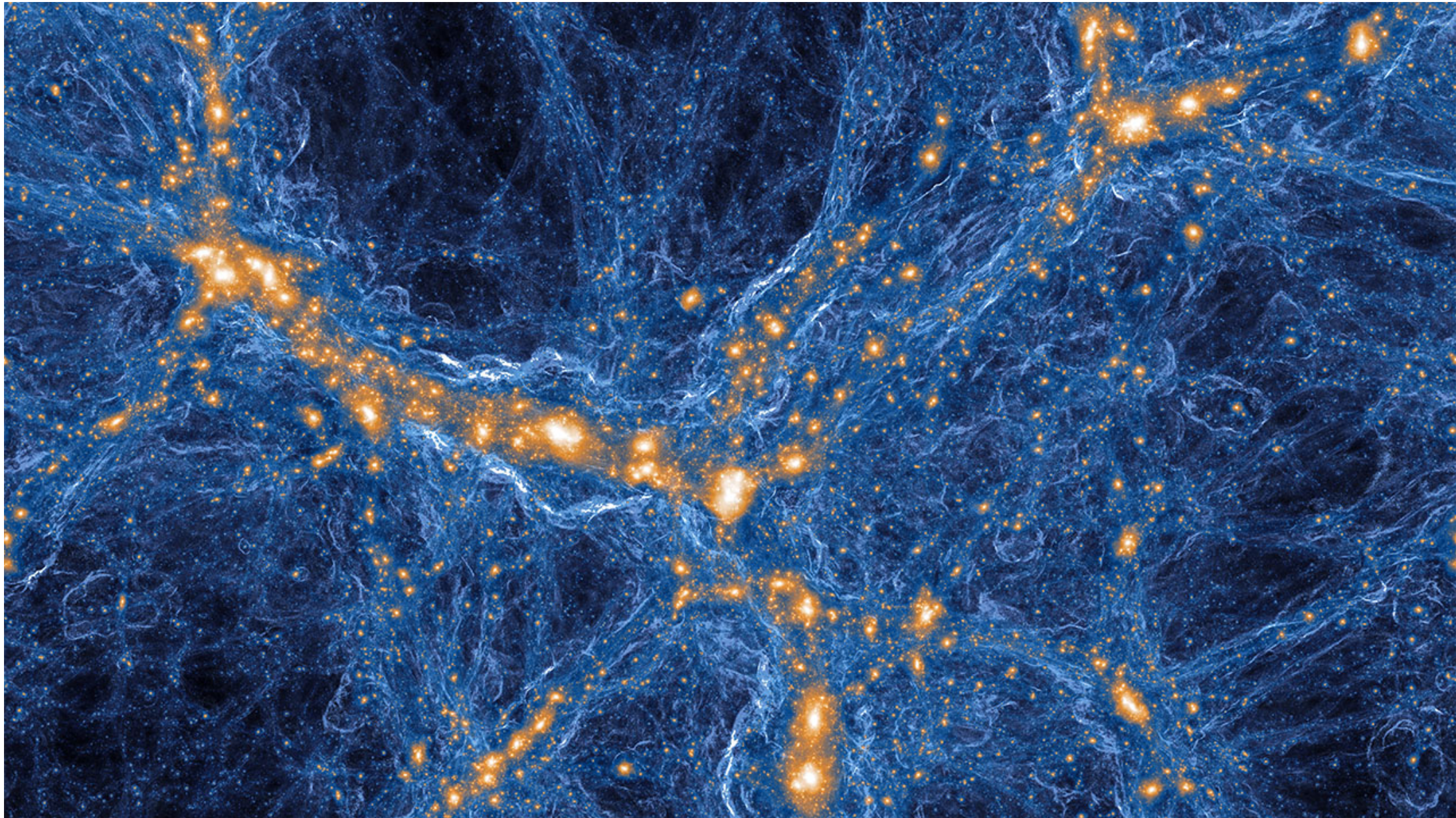
New physics  
(signal)

galaxy formation loops  
(LCDM background)





# Simplest example: galaxy bias



IllustrisTNG

$$\delta = \frac{\rho}{\langle \rho \rangle} - 1 \quad \delta_g = b_1 \delta_m + b_2 \delta_m^2 + \dots$$



$$\delta_g = b_1 \delta_m + b_2 \delta_m^2 + \dots$$



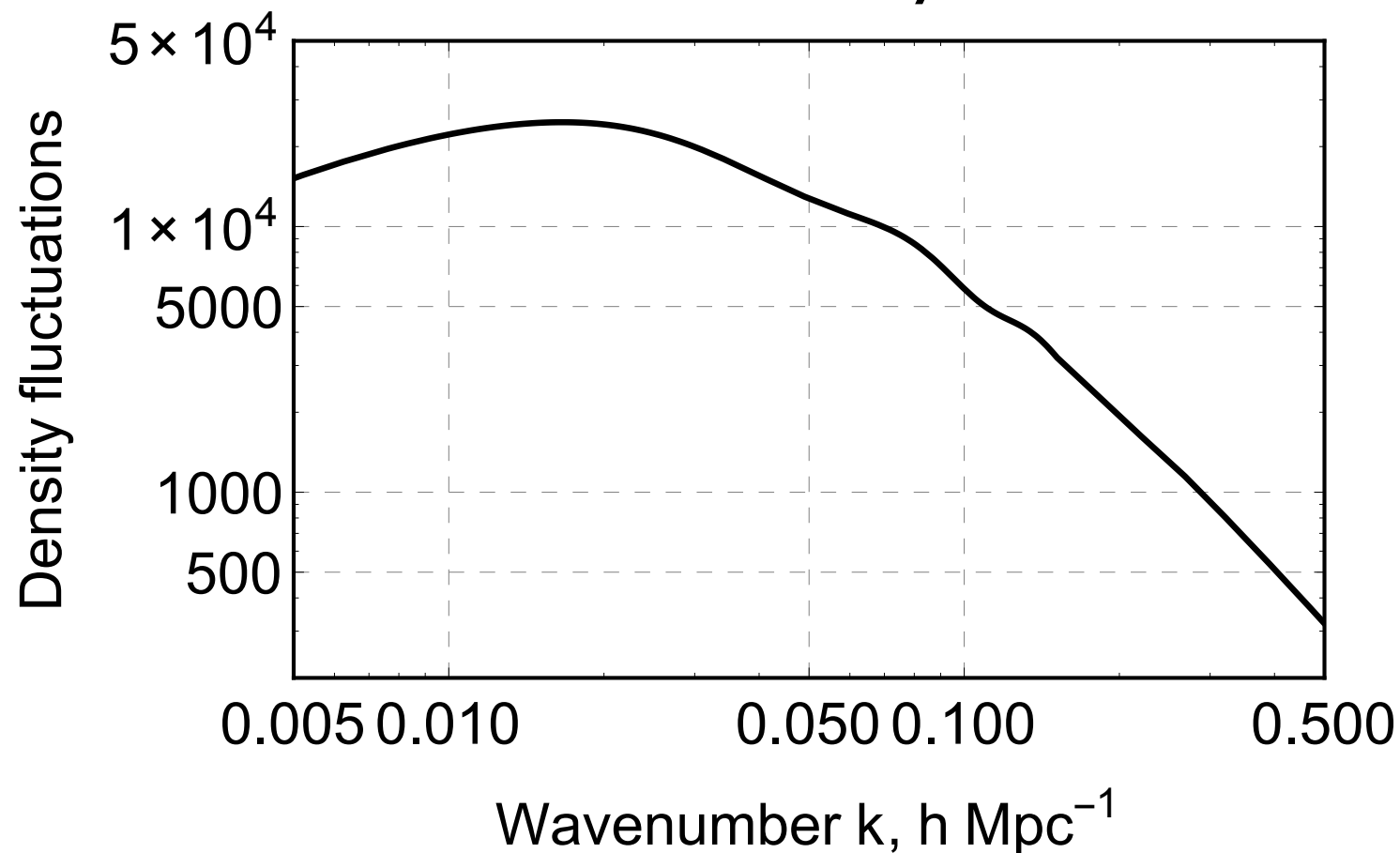
**Statistics:**

$$\langle \delta_m(\mathbf{k}) \delta_m(\mathbf{k}') \rangle = \delta_D^{(3)}(\mathbf{k} + \mathbf{k}') P_{\text{lin}}(k)$$

Gaussian random field

linear theory

~ propagator



depends on radiation, baryons, DM, etc. from Boltzmann codes

$$\delta_g = b_1 \delta_m + b_2 \delta_m^2 + \dots$$



**Statistics:**

$$\langle \delta_m(\mathbf{k}) \delta_m(\mathbf{k}') \rangle = \delta_D^{(3)}(\mathbf{k} + \mathbf{k}') P_{\text{lin}}(k)$$

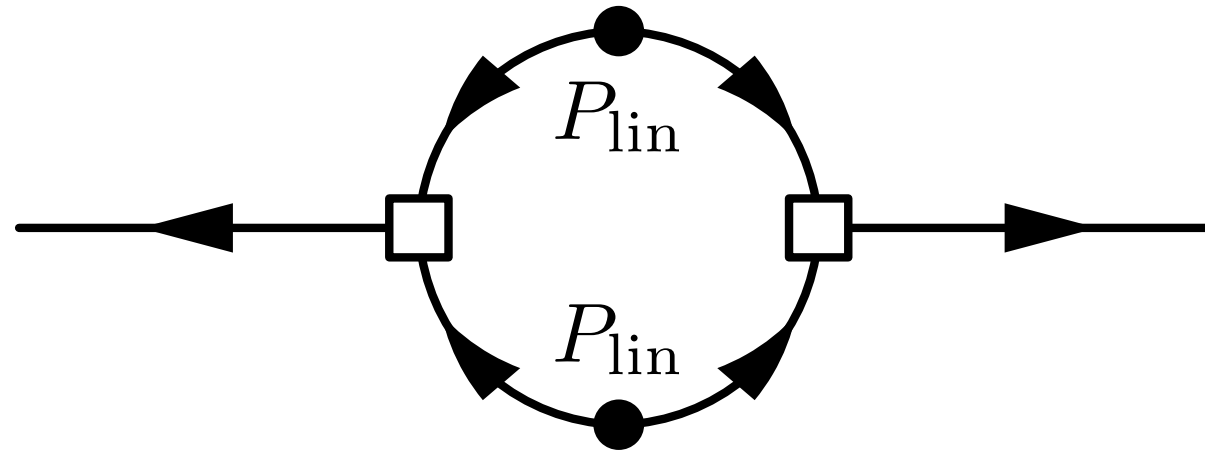
Non - linear theory

$$\langle \delta_g(\mathbf{k}) \delta_g(\mathbf{k}') \rangle = b_1^2 \underbrace{\langle \delta_m \delta_m \rangle}_{\text{bracket}} + 2b_1 b_2 \underbrace{\langle \delta_m \delta_m \delta_m \rangle}_{\text{bracket}} + b_2^2 \underbrace{\langle \delta_m \delta_m \delta_m \delta_m \rangle}_{\text{bracket}}$$

$$\langle \delta_m \delta_m \delta_m \delta_m \rangle' = \int \frac{d^3 q}{(2\pi)^3} P_{\text{lin}}(q) P_{\text{lin}}(|\mathbf{k} - \mathbf{q}|)$$

# Galaxy loops

$$\langle \delta_m \delta_m \delta_m \delta_m \rangle' = \int \frac{d^3 q}{(2\pi)^3} P_{\text{lin}}(q) P_{\text{lin}}(|\mathbf{k} - \mathbf{q}|)$$



MI “EFT for Large Scale Structure” ([2212.08488](#))

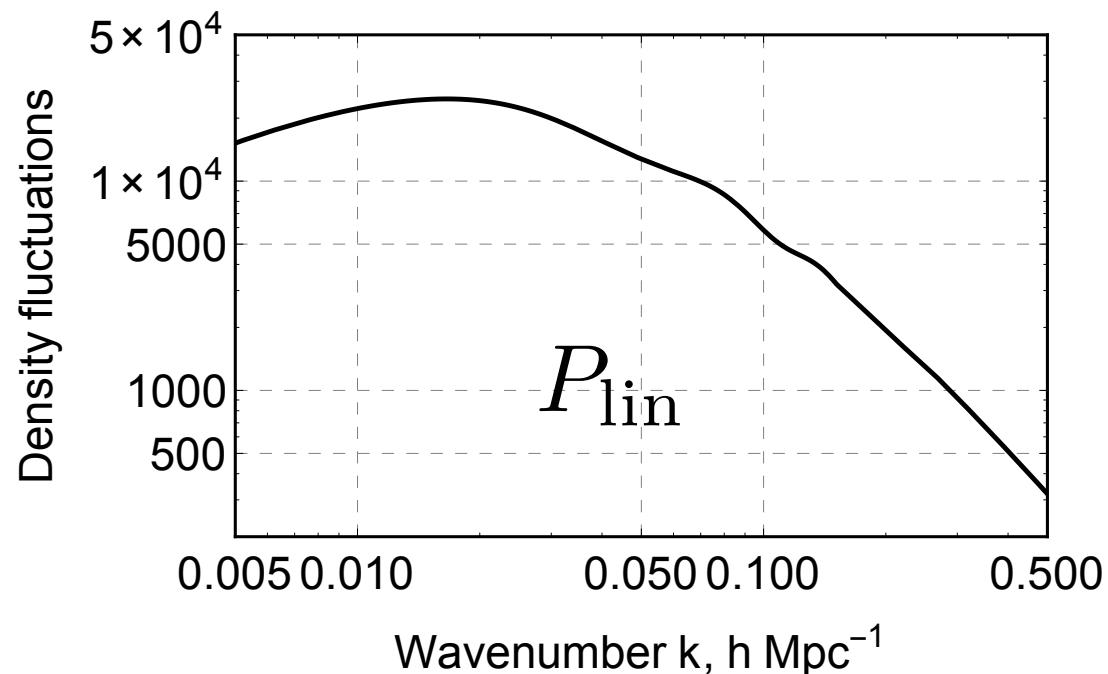
Baumann, Nicolis, Senatore, Zaldarriaga (2012)



# Loop evaluation: FFTLog

Simonovic, Zaldarriaga+ (2017)

$$\langle \delta_m \delta_m \delta_m \delta_m \rangle' = \int \frac{d^3 q}{(2\pi)^3} P_{\text{lin}}(q) P_{\text{lin}}(|\mathbf{k} - \mathbf{q}|)$$



$$P_{\text{lin}}(k) \approx \sum_m c_m k^{i\eta_m}$$

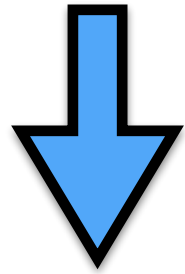
$$= \sum_m c_m e^{i\eta_m \ln k}$$

$$\int \frac{d^3 q}{(2\pi)^3} \frac{1}{q^{\nu_1} |\mathbf{k} - \mathbf{q}|^{\nu_2}} = \frac{k^{3-\nu_1-\nu_2}}{8\pi^{3/2}} \frac{\Gamma\left(\frac{3}{2} - \frac{\nu_1}{2}\right) \Gamma\left(\frac{3}{2} - \frac{\nu_2}{2}\right) \Gamma\left(\frac{\nu_1+\nu_2}{2} - \frac{3}{2}\right)}{\Gamma(\nu_1/2) \Gamma(\nu_2/2) \Gamma(3 - \nu_1/2 - \nu_2/2)}$$

Smirnov (1991)

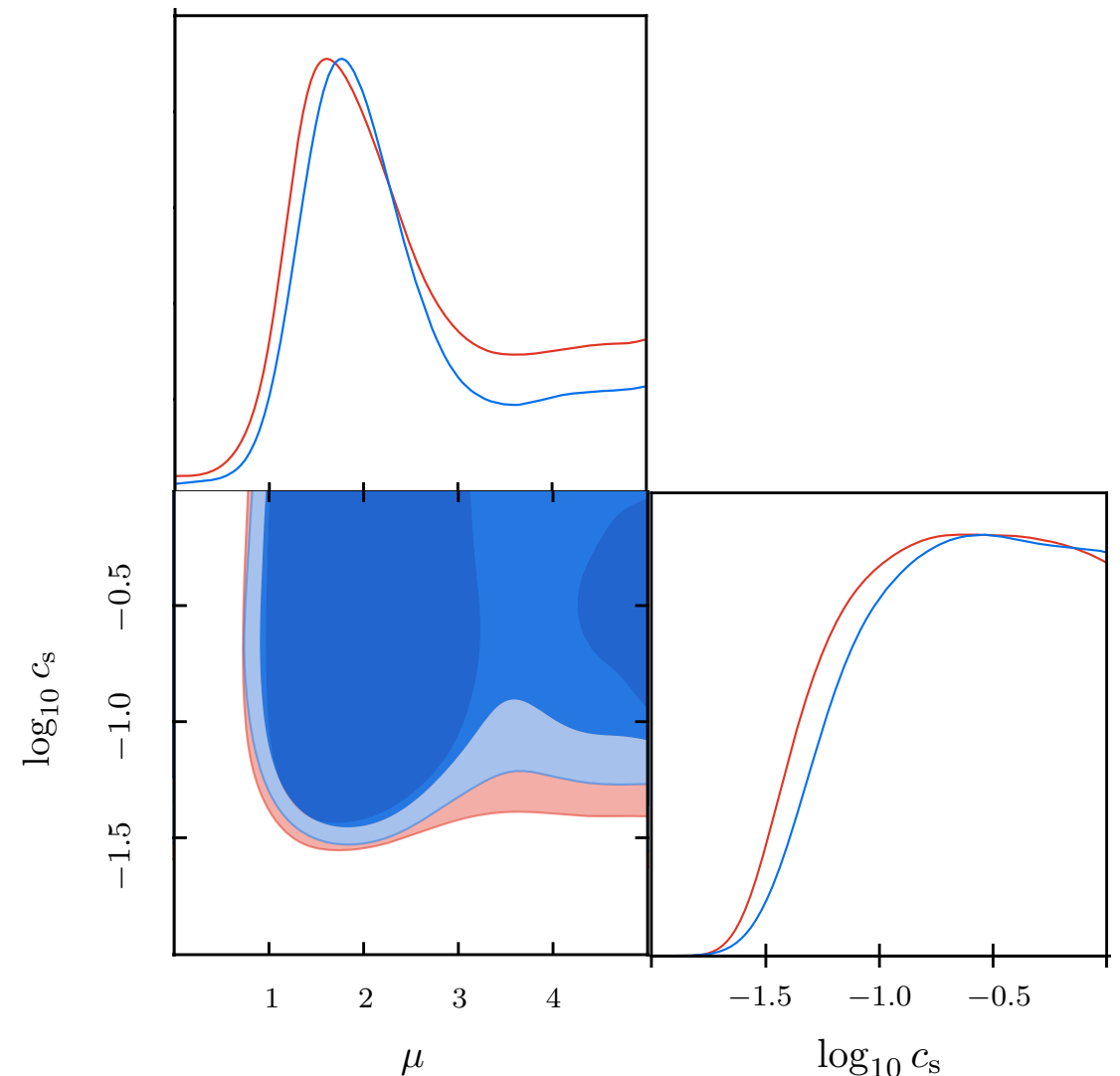
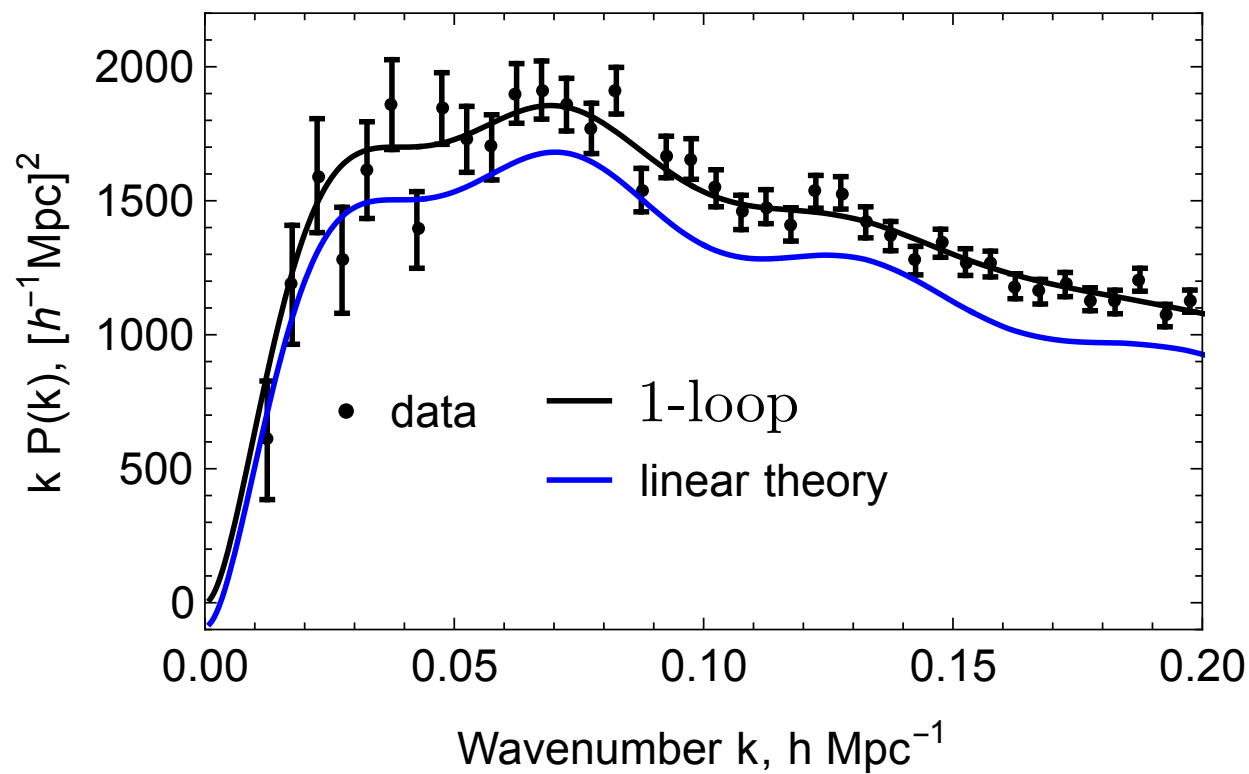
## 3D massless Euclidean QFT

$$\int \frac{d^3 q}{(2\pi)^3} \frac{1}{q^{\nu_1} |\mathbf{k} - \mathbf{q}|^{\nu_2}} = \frac{k^{3-\nu_1-\nu_2}}{8\pi^{3/2}} \frac{\Gamma\left(\frac{3}{2} - \frac{\nu_1}{2}\right) \Gamma\left(\frac{3}{2} - \frac{\nu_2}{2}\right) \Gamma\left(\frac{\nu_1+\nu_2}{2} - \frac{3}{2}\right)}{\Gamma(\nu_1/2) \Gamma(\nu_2/2) \Gamma(3 - \nu_1/2 - \nu_2/2)}$$



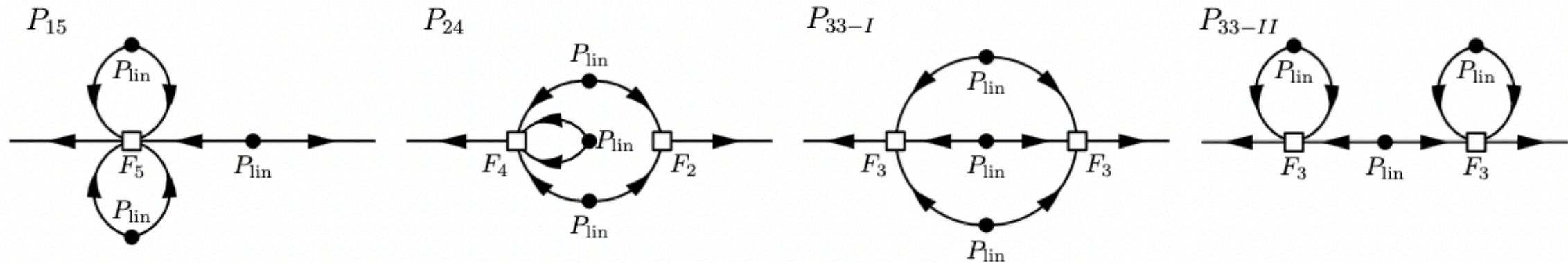
CLASS-PT code, 150 loops in 1 sec.

*Ivanov, Simonovic, Zaldarriaga (2019)*



# Two Loops

Simonovic, Zaldarriaga+ (2017)



$$\int_{\mathbf{q}} \frac{1}{q^{2\nu_4} |\mathbf{k} - \mathbf{q}|^{2\nu_5}} \int_{\mathbf{p}} \frac{1}{p^{2\nu_1} |\mathbf{k} - \mathbf{p}|^{2\nu_2} |\mathbf{q} - \mathbf{p}|^{2\nu_3}} \equiv k^{6-2\nu_{12345}} K(\nu_1, \dots, \nu_5)$$

$$K(\nu_1, \nu_2, \nu_3, \nu_4, \nu_5) = \frac{\sec(\pi\nu_{23})}{8\sqrt{\pi}\Gamma(\nu_1)\Gamma(\nu_2)\Gamma(\nu_3)\Gamma(3-\nu_{123})}$$

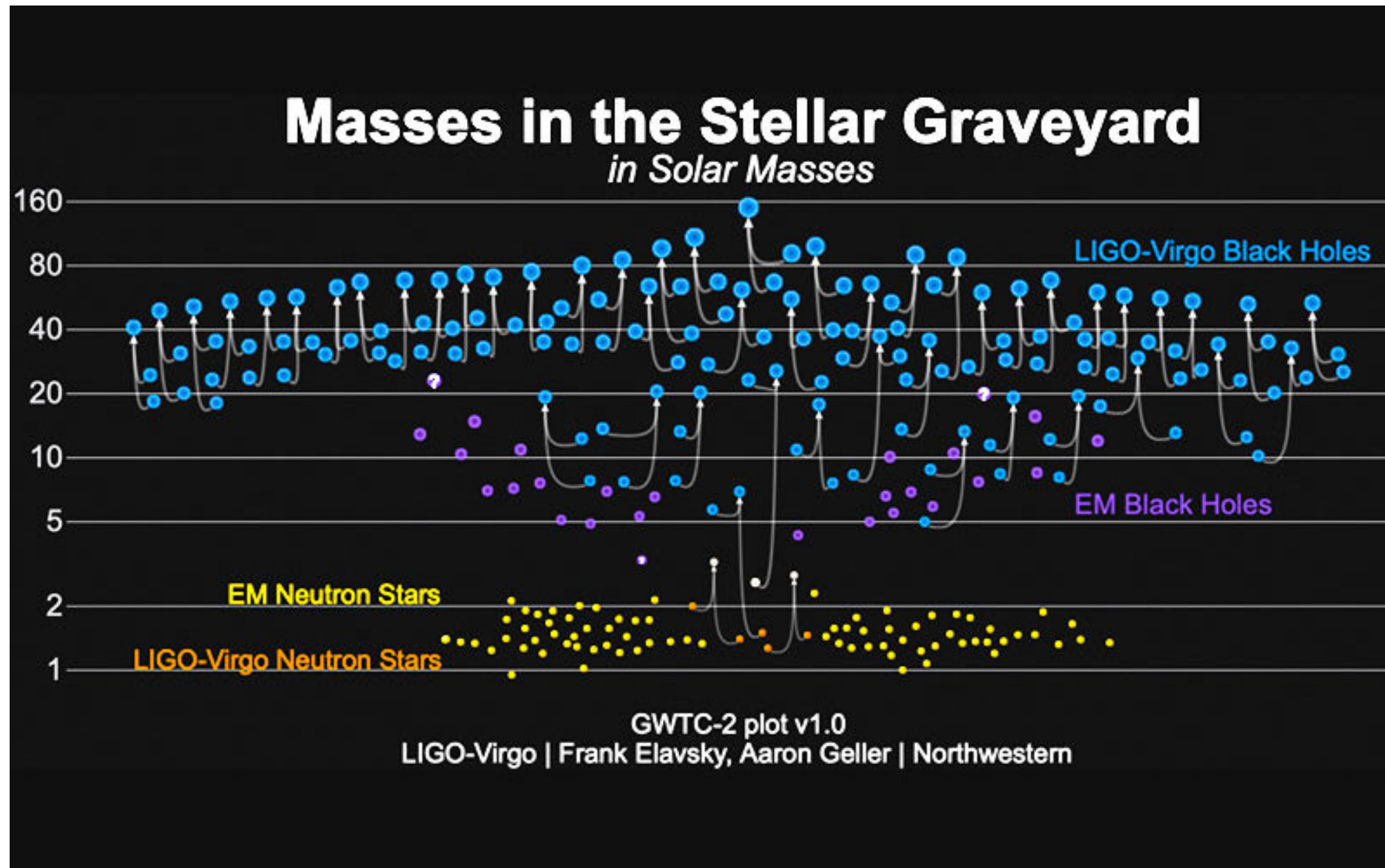
$$\times \sum_{n=0}^{\infty} [a_n(\nu_1, \nu_2, \nu_3) \kappa\left(\frac{3}{2} - \nu_{235} + n, -\nu_4, \nu_1 + n, \frac{3}{2} - \nu_2 + n, 3 - \nu_{23} + 2n\right) - a_n(\tilde{\nu}_1, \tilde{\nu}_2, \tilde{\nu}_3) \kappa\left(-\nu_5 + n, \frac{3}{2} - \nu_{134}, \tilde{\nu}_1 + n, \frac{3}{2} - \tilde{\nu}_2 + n, 3 - \tilde{\nu}_{23} + 2n\right)]$$

$$\kappa(\alpha, \beta, a, b, c) = \frac{1}{1+\alpha} \left[ \frac{1}{1+\beta} {}_3F_2\left(\begin{matrix} 1, a, b \\ 2 + \beta, c \end{matrix}; 1\right) - \dots \right]$$

- The FFTLog formula converges too slow
- New ideas from you are welcome!

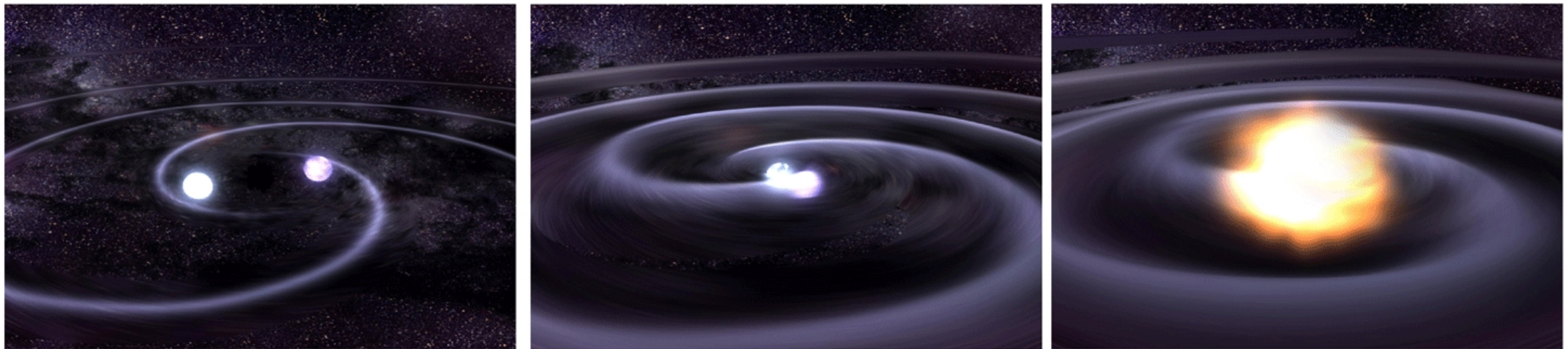


# The dawn of precision GW science



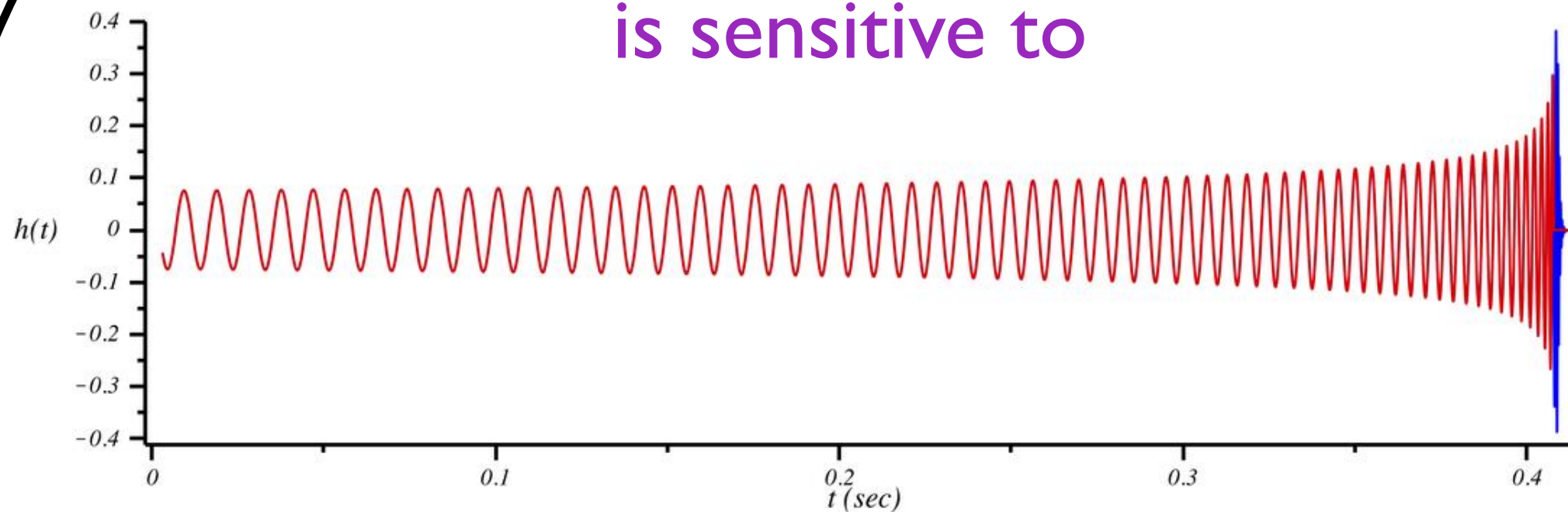
also see Zvi Bern's, Alessandra Buonanno's +++ talks

# A typical binary merger



© NASA

is sensitive to



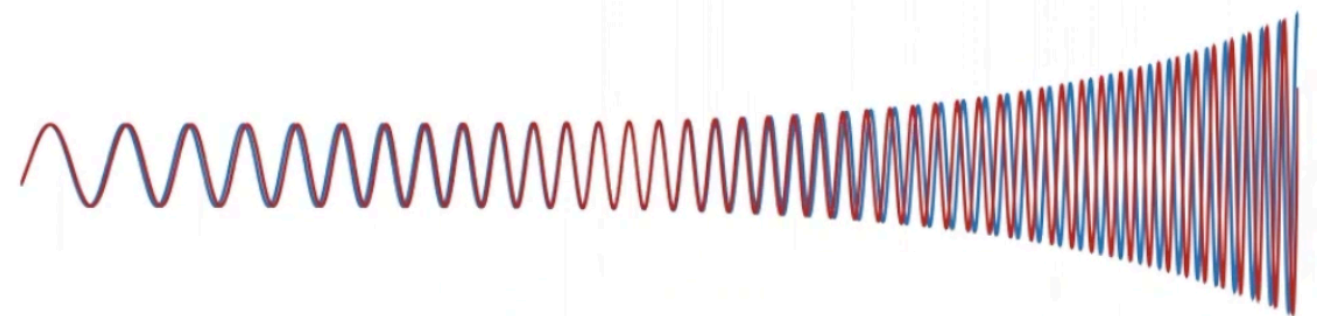
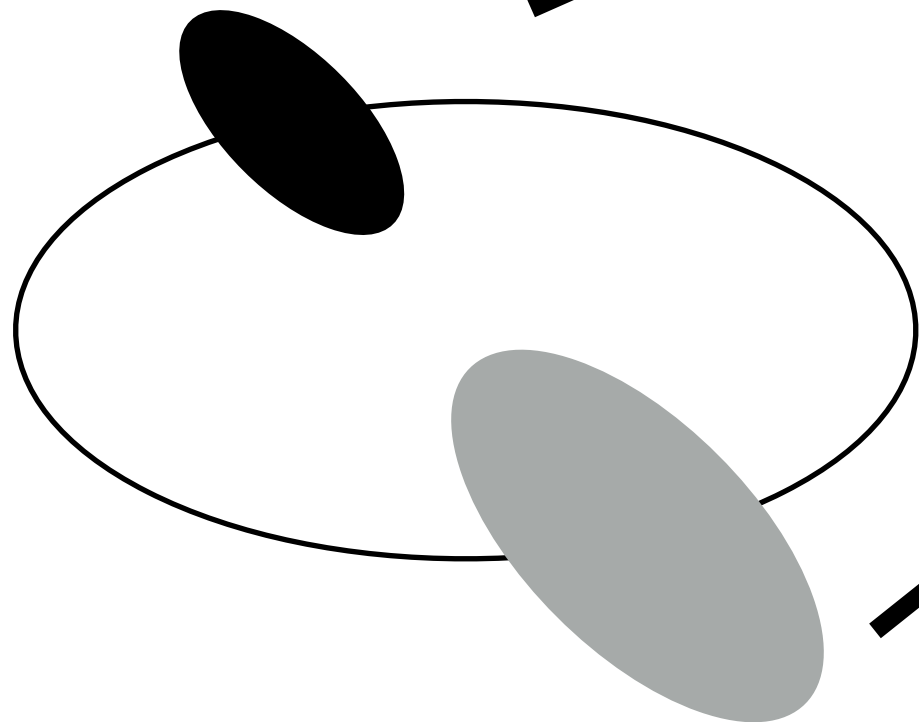
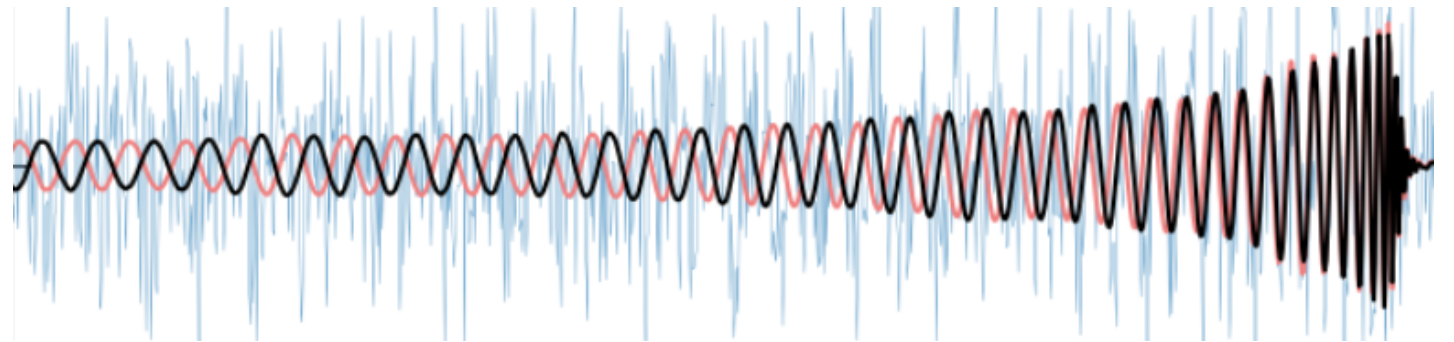
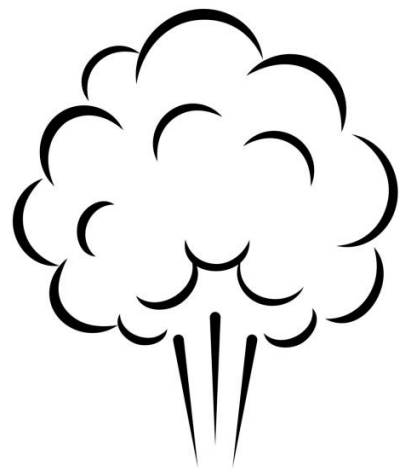
Quasinormal modes (QNMs)

tidal deformation of the sources



# Tides probe the nature of compact objects

heating (4 PN),  
dissipative



deformation (5 PN),  
conservative Love numbers



One can test EoS for neutron stars

*Flanagan, Hinderer (2007) ++*

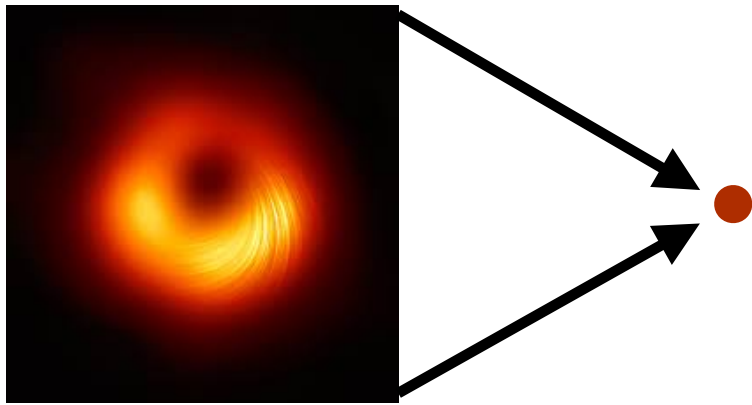


Use LVK data to test exotic compact objects

*Chia, Zhou, MI (2404.14641)*



# Love numbers in Worldline EFT



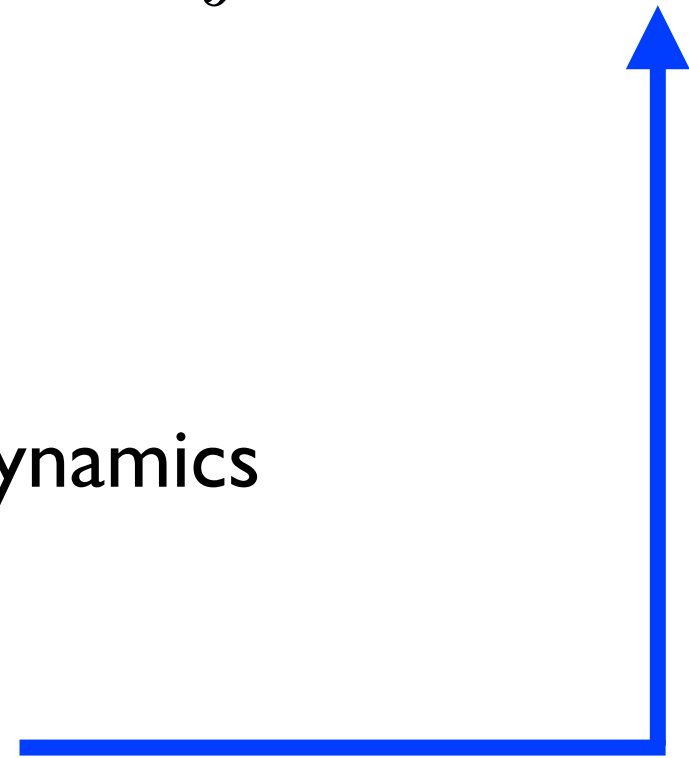
at  
 $r \gg r_s$

a black hole is described  
 by the world line  
 effective action

$$S_{\text{eff}} = -m \int ds + \lambda_2 R^5 \int ds E_{\mu\nu} E^{\mu\nu} + \int ds E_{\mu\nu} Q^{\mu\nu}(X)$$

$$E_{ij} = \partial_i \partial_j \Phi \quad \leftarrow \quad u^\mu u^\lambda R_{\mu\nu\lambda\rho}$$

- Local operators capture conservative dynamics
- Wilson coefficients = Love numbers
- 2pf of Q encodes dissipative dynamics



# Matching of Tides in Worldline EFT

$$S_{\text{eff}} = -m \int ds + \lambda_2 R^5 \int ds E_{\mu\nu} E^{\mu\nu} + \int ds E_{\mu\nu} Q^{\mu\nu}(X)$$



Wilson coeffs extracted from GR matching calculations, use for new predictions



Normally done with off-shell quantities (~Newton potential)



Works well\* @LO (static LNs), but matching conditions are trivial

*Damour, Poisson, +’09+*



@NLO (dynamical LNs) some confusion: gauge, coordinate dependence

$$\int ds \dot{E}_{\mu\nu} \dot{E}^{\mu\nu}$$

e.g. *Poisson (2021a,b)* vs *Chakrabarti +’13, Charalambous, Dubovsky, MI’21*

*Saketh, Zhou, MI’21*



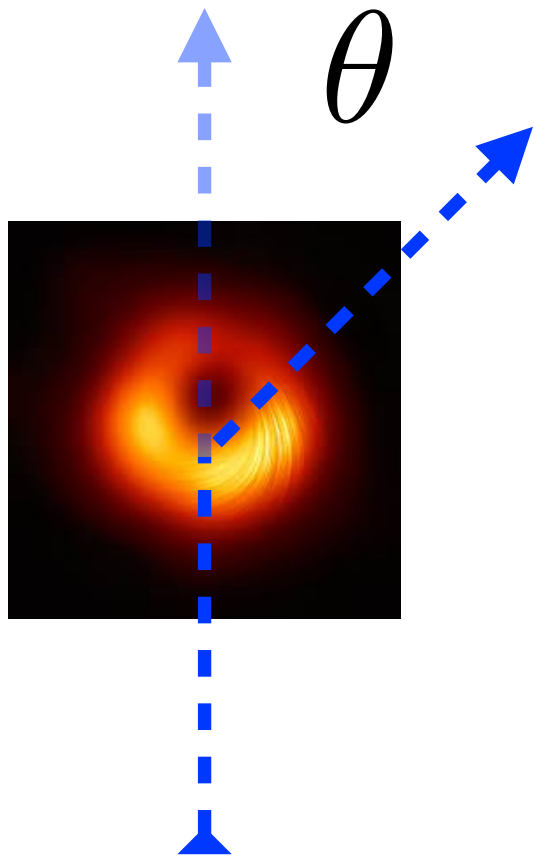
Alternative: on-shell scattering amplitudes

Free of gauge, coordinate, and field-redefinition

*MI, Li, Parra-Martinez, Zhou’24*

# Gravitational Raman Scattering

MI, Li, Parra-Martinez, Zhou'24



$\phi, A^\mu, h_{\mu\nu}^{TT}$

frequency  $\omega$

$$S_{\text{eff}} = -m \int ds + \int ds [C_1 R^3 (\partial_i \phi)^2 + C_{0,\omega^2} R^3 \dot{\phi}^2 + C_{1,\omega^2} R^5 (\partial_i \dot{\phi})^2] + \int ds [\phi Q(X) + \partial_i \phi Q^i(X) + \dots]$$



Partial wave basis:

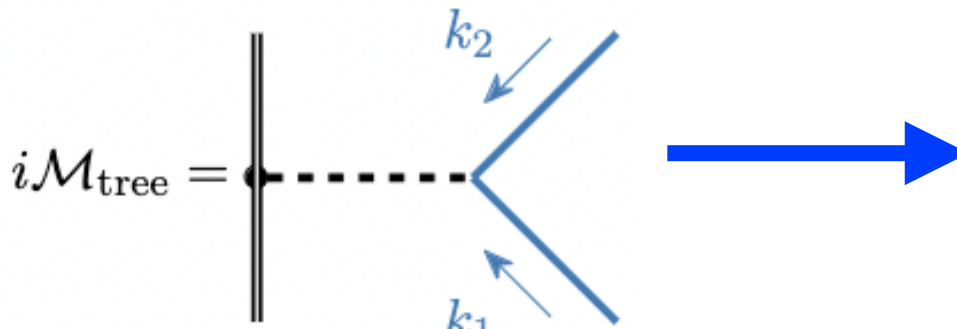
$$i\mathcal{M}(\omega, \theta) = \frac{2\pi}{\omega} \sum_{\ell=0}^{\infty} (2\ell + 1) (\eta_\ell e^{2i\delta_\ell} - 1) P_\ell(\cos \theta)$$

 (Real part of) phase encode conservative tides (Love numbers)

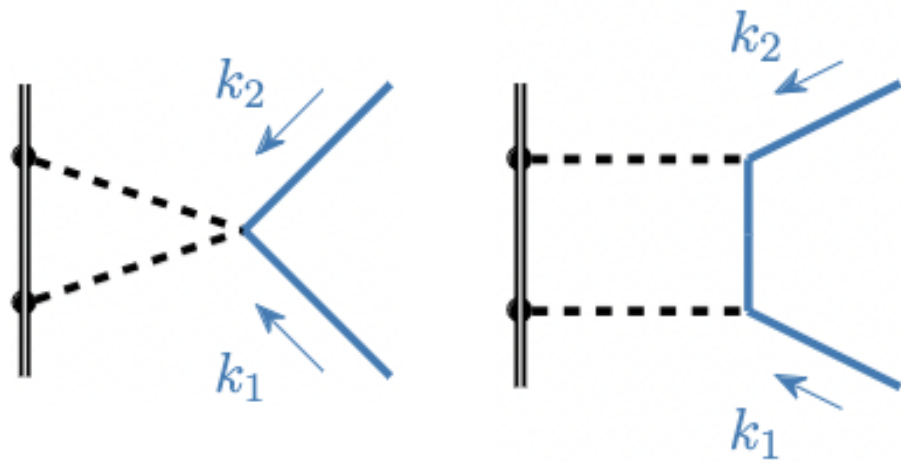
 Inelasticity encodes tidal heating

# Non-linearity of Gravity

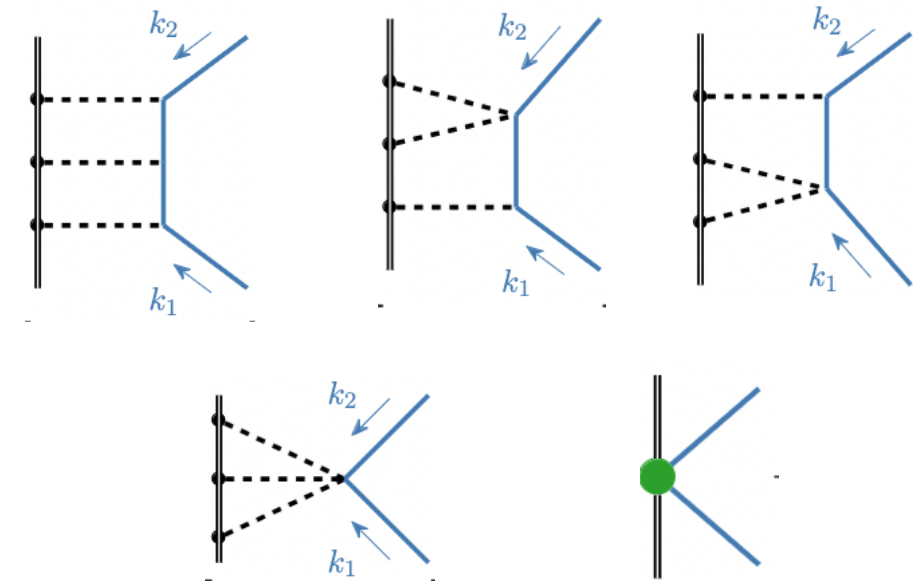
$$i\mathcal{M}(\omega) \sim \underbrace{(mG)}_{\text{“Rutherford”}} \underbrace{(1 + mG\omega + (mG\omega)^2 + \dots)}_{\text{relativ. corrections}} + \underbrace{(mG)^5 \omega^4}_{\text{tides}}$$

**1PM:**  $i\mathcal{M}_{\text{tree}} =$    $\frac{d\sigma}{d\Omega} \Big|_{1\text{PM}} = \frac{G^2 m^2}{\sin^2\left(\frac{\theta}{2}\right)} \cos^4\left(\frac{\theta}{2}\right)$

**2PM:**



**3PM:**




3PM calculated for the first time thanks to IBP+diff. equations

Background field method integrands [Cheung, Parra-Martinez, Rothstein+'23](#)



## Results:


**3PM EFT:**

$$\delta_\ell \Big|_{\text{EFT}} = -\frac{\lambda}{2\epsilon_{\text{IR}}} + \frac{\lambda}{2} \ln \left( \frac{4\omega^2}{\bar{\mu}^2} \right) + \sum_{n=1}^3 \nu_n^\ell \lambda^n + \delta_\ell^{G^3},$$

$$1 - \eta_\ell \Big|_{\text{EFT}} = \frac{\ell! \omega^{2\ell+1} \text{Im} F_\ell(\omega)}{2\pi(2\ell+1)!!} \left( 1 + \pi\lambda + \lambda^2 \eta_\ell^{G^2} \right), \quad \lambda \equiv 2Gm\omega \ll 1$$

$$\delta_0^{G^3} \Big|_{\text{EFT}} = \lambda^3 \left[ \frac{1}{4\epsilon_{\text{UV}}} + \frac{49}{24} - \frac{1}{2} \ln \left( \frac{4\omega^2}{\bar{\mu}^2} \right) \right] + \frac{C_{0,\omega^2}\omega^3}{4\pi}, \quad \eta_0^{G^2} \Big|_{\text{EFT}} = \frac{67}{12} - \frac{11}{6} \left( -\frac{1}{2\epsilon_{\text{UV}}} + \ln \left( \frac{4\omega^2}{\bar{\mu}^2} \right) \right) + \frac{\pi^2}{3}$$

$$\delta_1^{G^3} \Big|_{\text{EFT}} = \frac{C_1\omega^3}{12\pi} \left( 1 + \pi\lambda + \lambda^2 \eta_1^{G^2} \right) + \frac{C_{1,\omega^2}\omega^5}{12\pi}, \quad \eta_1^{G^2} \Big|_{\text{EFT}} = \frac{413}{100} - \frac{19}{30} \left( -\frac{1}{2\epsilon_{\text{UV}}} + \ln \left( \frac{4\omega^2}{\bar{\mu}^2} \right) \right) + \frac{\pi^2}{3}$$


**Comparison with known results for black holes in GR:**

$$\delta_\ell \Big|_{\text{GR}} = (r_s \omega) \ln(2\omega r_s) + \sum_{n=1}^3 \nu_n^\ell (r_s \omega)^n + \delta_\ell^{G^3} \quad \delta_0^{G^3} \Big|_{\text{GR}} = (r_s \omega)^3 \left[ \frac{7}{12} - \gamma_E - \ln(2r_s \omega) \right]$$

Static Love numbers are zero  
and do not run:  $C_1 = 0$

Dynamical Love numbers non zero  
and do run:

$$C_{0,\omega^2}(\bar{\mu})^{\overline{\text{MS}}} = -4\pi r_s^3 \left[ \frac{1}{4\epsilon_{\text{UV}}} + \ln(\bar{\mu} r_s) + \frac{35}{24} + \gamma_E \right]$$

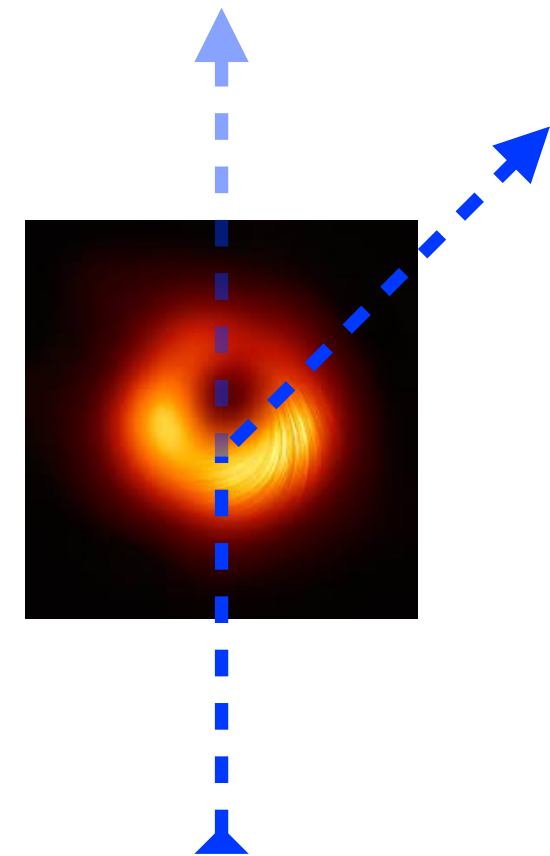
# Scalar Raman Scattering: Upshot

- Perfect match: consistency of the EFT and the renormalization program!
- IR/ UV logs neatly separated in EFT
- Photon/graviton results underway!
- Running part can be matched thanks to the factorization of GR amplitudes

*Chakrabarti ++'13, Charalambous, Dubovsky, MI'21  
Zhou MI'22, Saketh, Zhou, MI'23*

- Dynamical Love numbers of BHs are non-zero & run,  
~ natural from the EFT perspective

Universal running: same for BH & NS



$$\int ds \dot{E}_{\mu\nu} \dot{E}^{\mu\nu}$$

# Summary



Large scale structure - a powerful probe of physics



Our Universe is the largest collider



Loops thanks to QFT, but need more :)



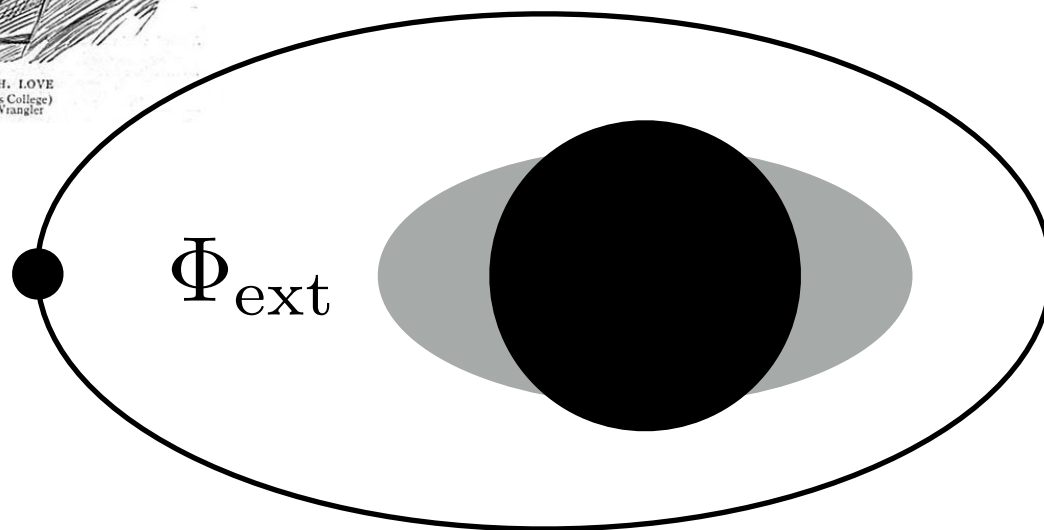
Gravitational waves - probe nature of compact objects



Tides understood thanks to scattering amplitudes

Thank you!

# Love numbers in Newtonian gravity



$$\Phi = \Phi_0 + \Phi_2 + \Phi_3 + \dots$$

$$\Phi_0 = -\frac{M}{r}$$

$$\Phi_2 = r^2 E_{ij} n^i n^j + \frac{Q_{ij} n^i n^j}{r^3}$$

tidal source      tidal response

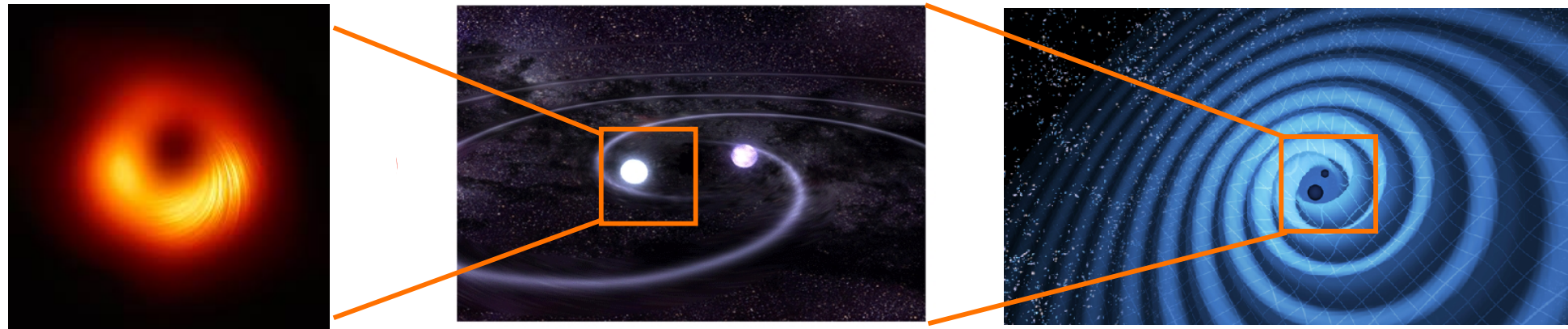
$$\Phi_\ell = r^\ell E_{i_1 \dots i_\ell} n^{i_1} \dots n^{i_\ell} + Q_{i_1 \dots i_\ell} \frac{n^{i_1} \dots n^{i_\ell}}{r^{\ell+1}}$$

Linear response theory:  $Q_\ell = \Lambda_\ell r_s^{2\ell+1} E_\ell$

**Love number**

$(G = c = \hbar = 1)$

# Worldline EFT



© LIGO collaboration, [ligo.org](http://ligo.org)

EFT framework: disentangle physics from different scales

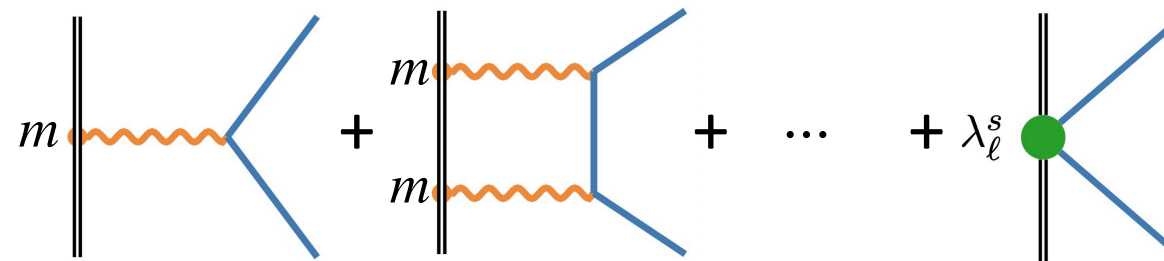
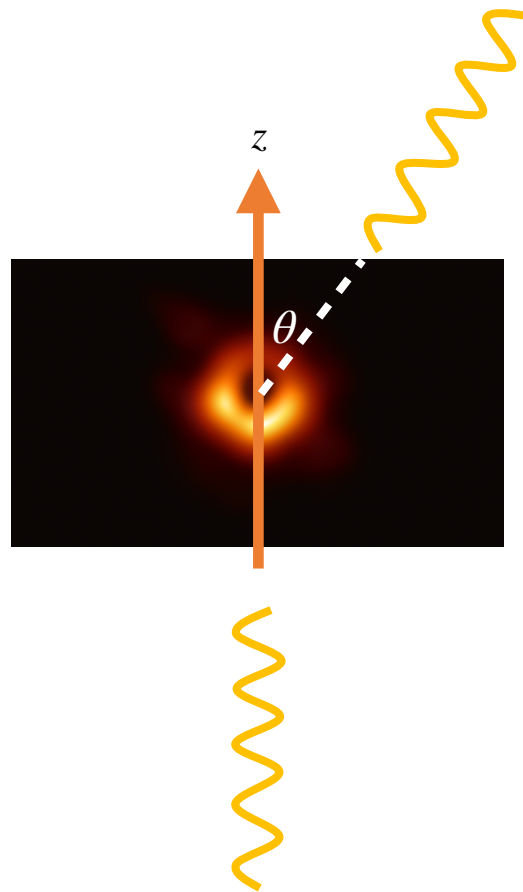


at  $r \gg r_s$  a black hole is described by the world line effective action

$$S = -m \int ds + \dots$$



# Matching LN via Raman scattering amplitudes



PM expansion:  $G + G^2 + \dots$

$$\lambda_\ell^s \sim R^{2\ell+1} \propto G^{\frac{2\ell}{D-3}+1}$$

Analytic continuation in angular momentum (Gribov-Froissart) separates LN and PM

$$iA \sim (Gm) + (Gm)^2 \omega + \dots + (Gm)^n \omega^{n-1} + (Gm)^{2\ell+1} \omega^{2\ell} + \dots$$

N.B. the limit  $\ell \rightarrow n \in \mathbb{N}$

leads to singularities if LNs run  
(there's an actual mixing between LN  
and PM corrections)

Zhou, MI'22

MI, Li, Parra-Martinez, Zhou'24

# Dynamical Love numbers

- Problem: RG running of LNs (PN loops + local counterterms)
- EFT gives a consistent definition via matching ( $\sim$ coupling running)
- Dynamical LNs of BHs are not zero and run *Chakrabarti ++'13, MI ++'21*

$$C_{\dot{E}^2} \int d\tau \dot{E}^{ij} \dot{E}_{ij} \quad C_{E\dot{E}} \int d\tau \varepsilon^{ijk} \hat{S}_i E_{jl} \dot{E}_k^l$$

- Running can be matched analytically with FG, e.g.

$$\frac{dC_{\dot{E}^2}}{d \log \mu} = \frac{32}{45} m^7 G^6 \quad \frac{dC_{E\dot{E}}}{d \log \mu} = -\frac{32}{45} \frac{a}{m} m^6 G^5$$

*Saketh, Zhou, MI' 23*

- Or extracted fully from Raman process amplitudes:

$$C_{0,\omega^2}(\bar{\mu})^{\overline{\text{MS}}} = -4\pi r_s^3 \left[ \frac{1}{4\epsilon_{\text{UV}}} + \ln(\bar{\mu} r_s) + \frac{35}{24} + \gamma_E \right] \sim \int \dot{\phi}^2$$

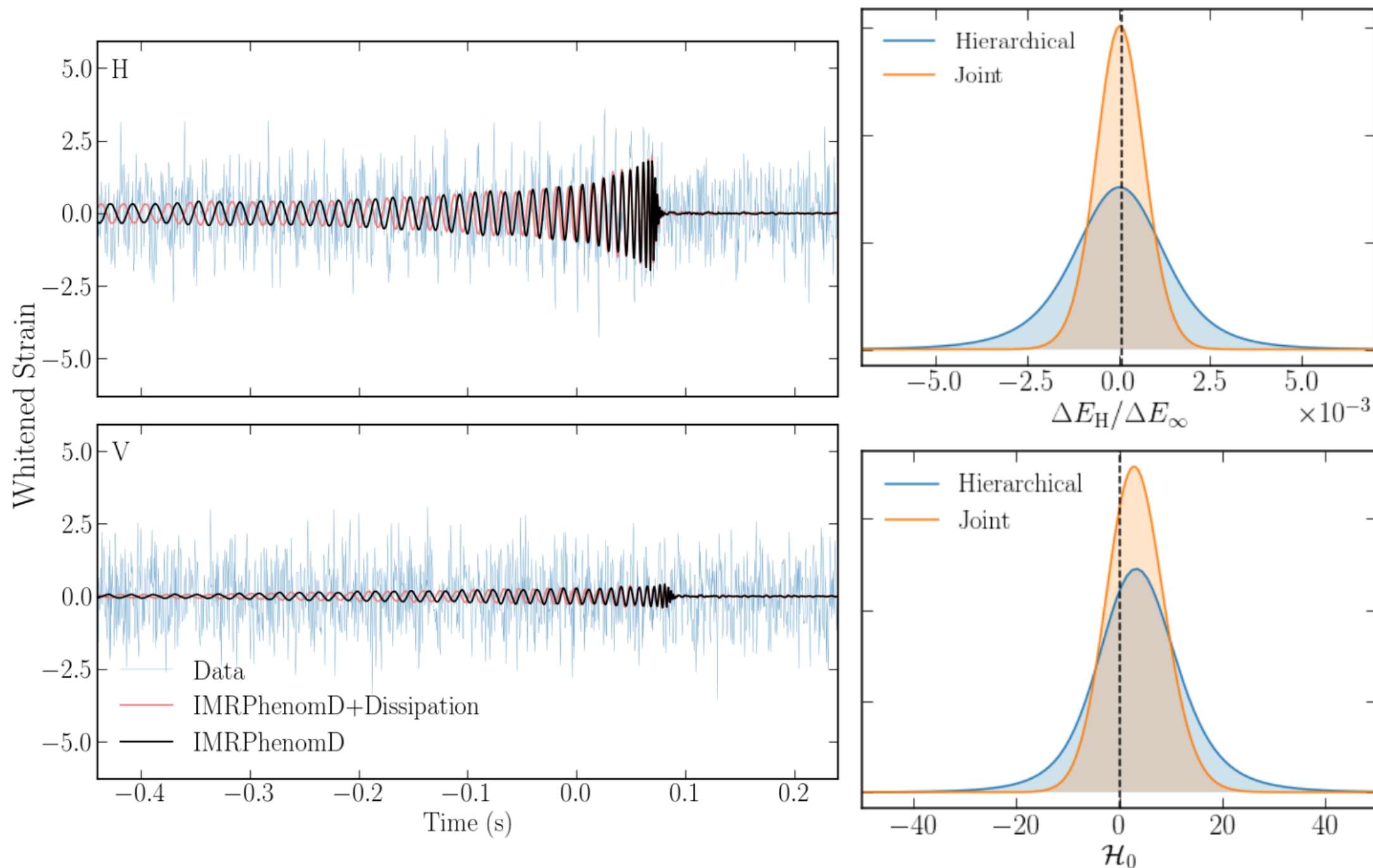
*MI, Li, Parra-Martinez, Zhou'24*



Worldline EFT: consistent definition of tidal effects in GR

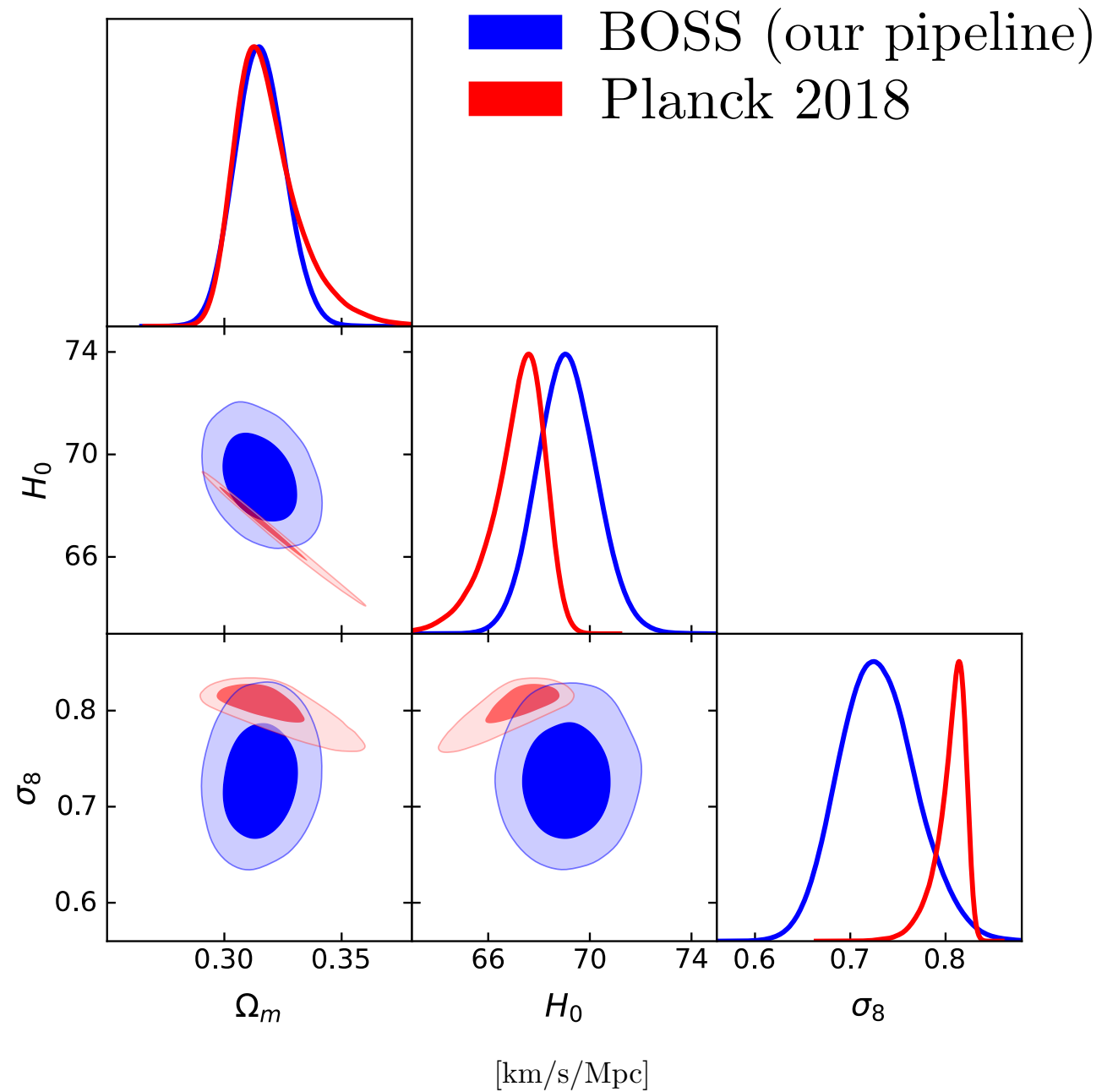
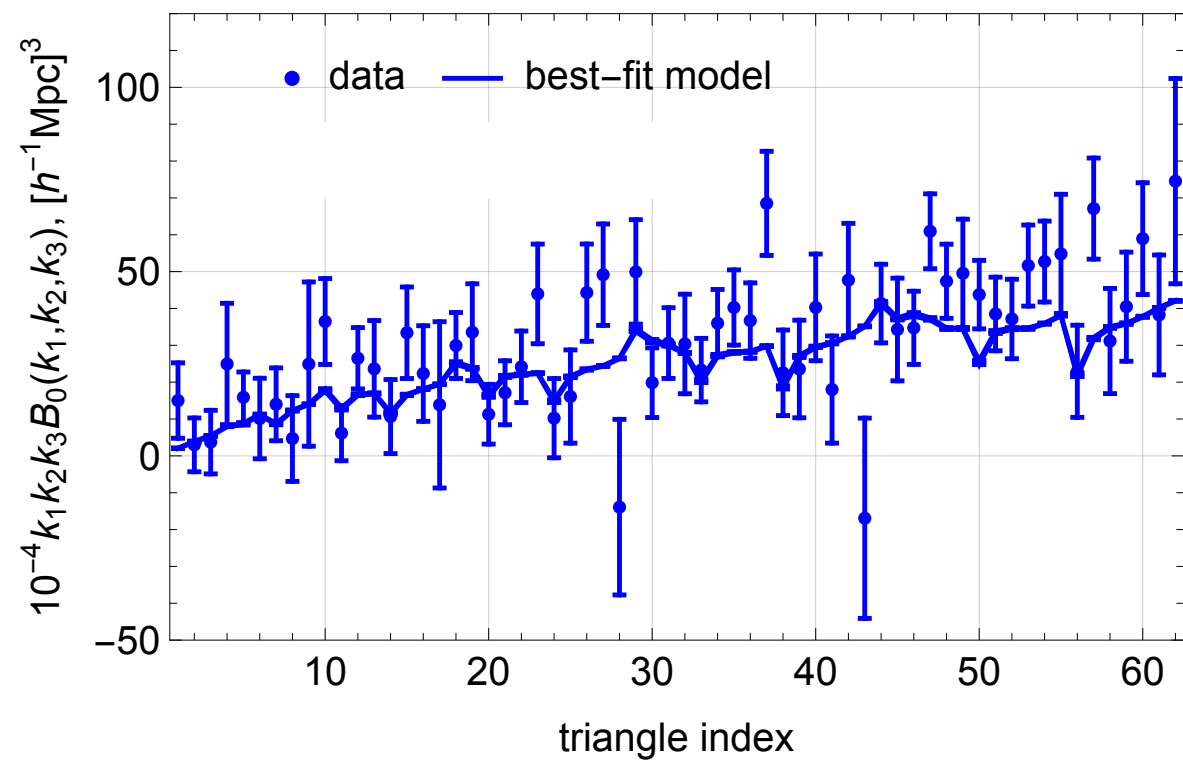
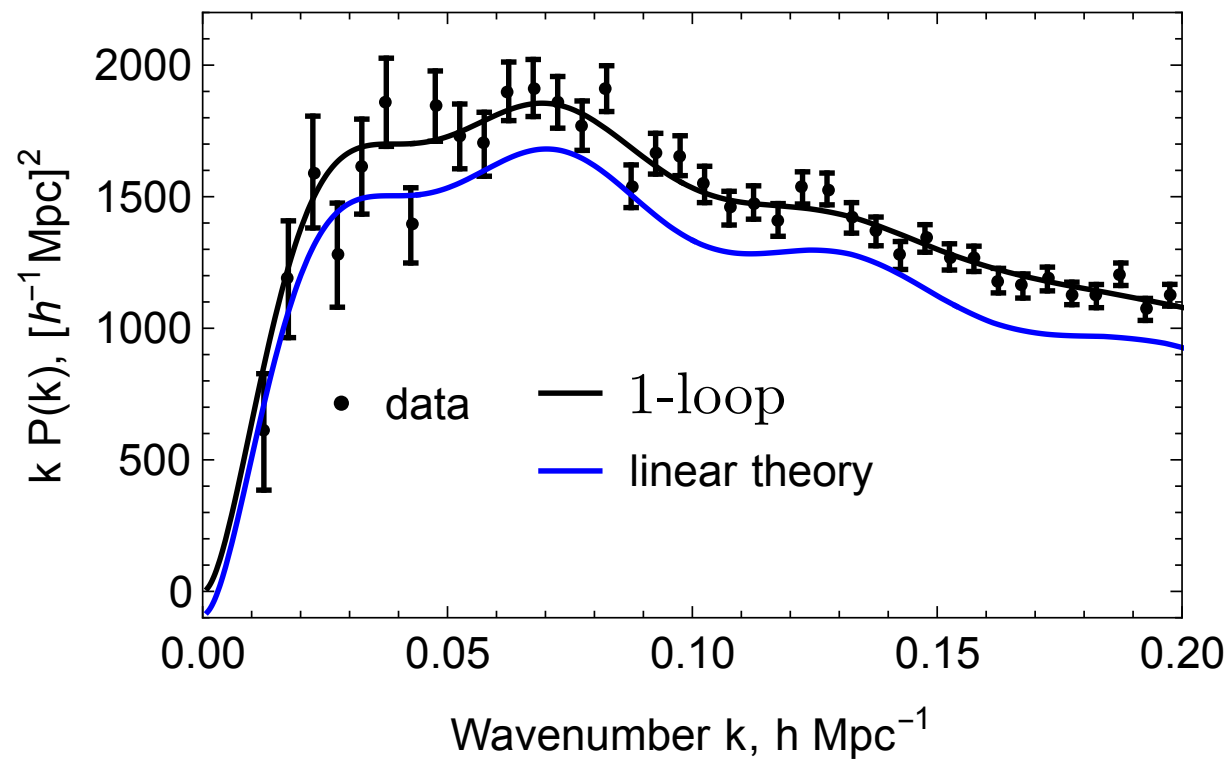


First constrains on exotic compact objects with LIGO +VK





# Large-scale structure: re-analysis of BOSS data



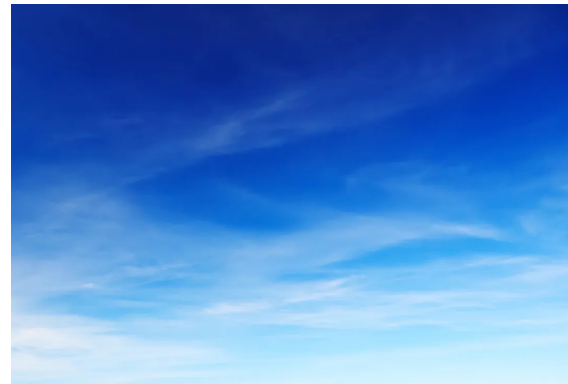
*MI, Simonovic, Zaldarriaga (2019), Philcox, MI (2021) ++  
D'Amico, Kokron++(2019), Chen, White, Vlah (2021)*

# Why is the sky blue ?

Scattering of light by neutral atoms (~ nitrogen in the air):

$$S_{\text{eff}} = Q \int \cancel{ds} A_0 + \chi \int ds E_i E^i + \dots \quad \chi \sim [\text{cm}]^3$$

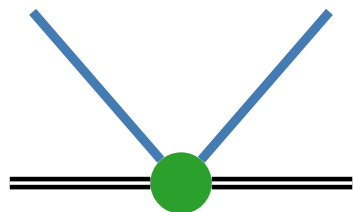
$$i\mathcal{M} \sim \chi \omega^2 \quad \longrightarrow$$



Scattering of GWs by a compact object:

$$S_{\text{eff}} = -m \int ds + \lambda_2 R^5 \int ds E_{\mu\nu} E^{\mu\nu}$$

$$i\mathcal{M} = \text{Diagram} \sim \lambda_2 R^5 \omega^4$$



Goldberger, Rothstein (2004,2007)  
Ivanov, Zhou (2022)

Universality of EFT dictates that the same param's appear in waveforms !

 Comparison with known results for black holes in GR:

$$\delta_\ell \Big|_{\text{GR}} = (r_s \omega) \ln(2\omega r_s) + \sum_{n=1}^3 \nu_n^\ell (r_s \omega)^n + \delta_\ell^{G^3} \quad \delta_0^{G^3} \Big|_{\text{GR}} = (r_s \omega)^3 \left[ \frac{7}{12} - \gamma_E - \ln(2r_s \omega) \right]$$

Renormalization flow of dynamical tides:

$$C_{0,\omega^2}(\bar{\mu})^{\overline{\text{MS}}} = -4\pi r_s^3 \left[ \frac{1}{4\epsilon_{\text{UV}}} + \ln(\bar{\mu} r_s) + \frac{35}{24} + \gamma_E \right]$$

$$\text{Im} F_0(\omega; \bar{\mu})^{\overline{\text{MS}}} = 4\pi r_s^2 |\omega| \left( 1 + (r_s \omega)^2 \left[ \frac{\pi^2}{3} - \frac{5}{18} - \frac{11}{3} \left( \ln(\bar{\mu} r_s) + \gamma_E \right) \right] \right) .$$

Static Love numbers are zero and do not run:  $C_1 = 0$

 Perfect match: consistency of the EFT and the renormalization program!



# Loops in Large Scale Structure

● Perturbation theory + Ensemble average = Loops

● Particle physics: ensemble of virtual particles

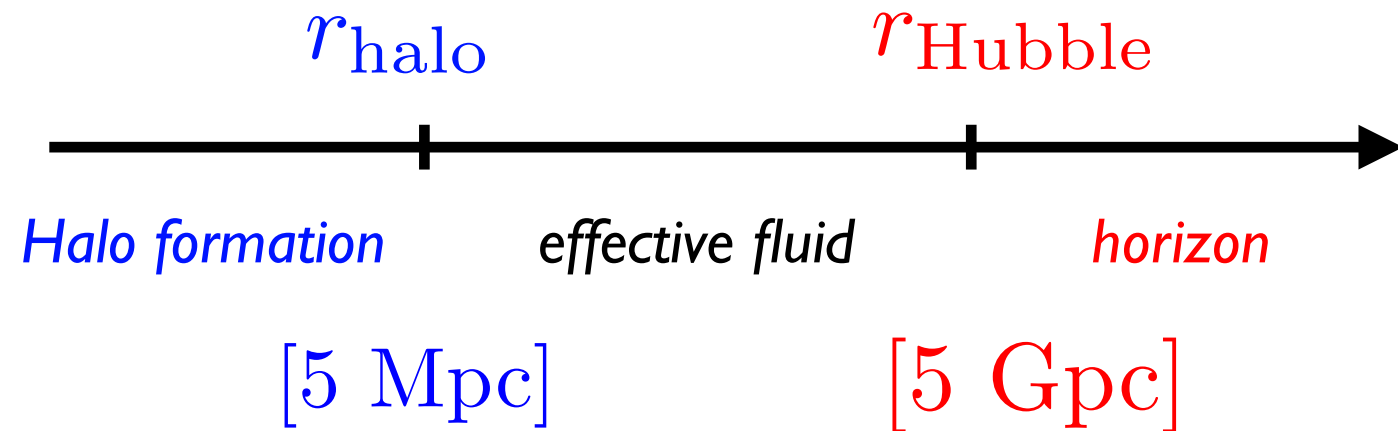
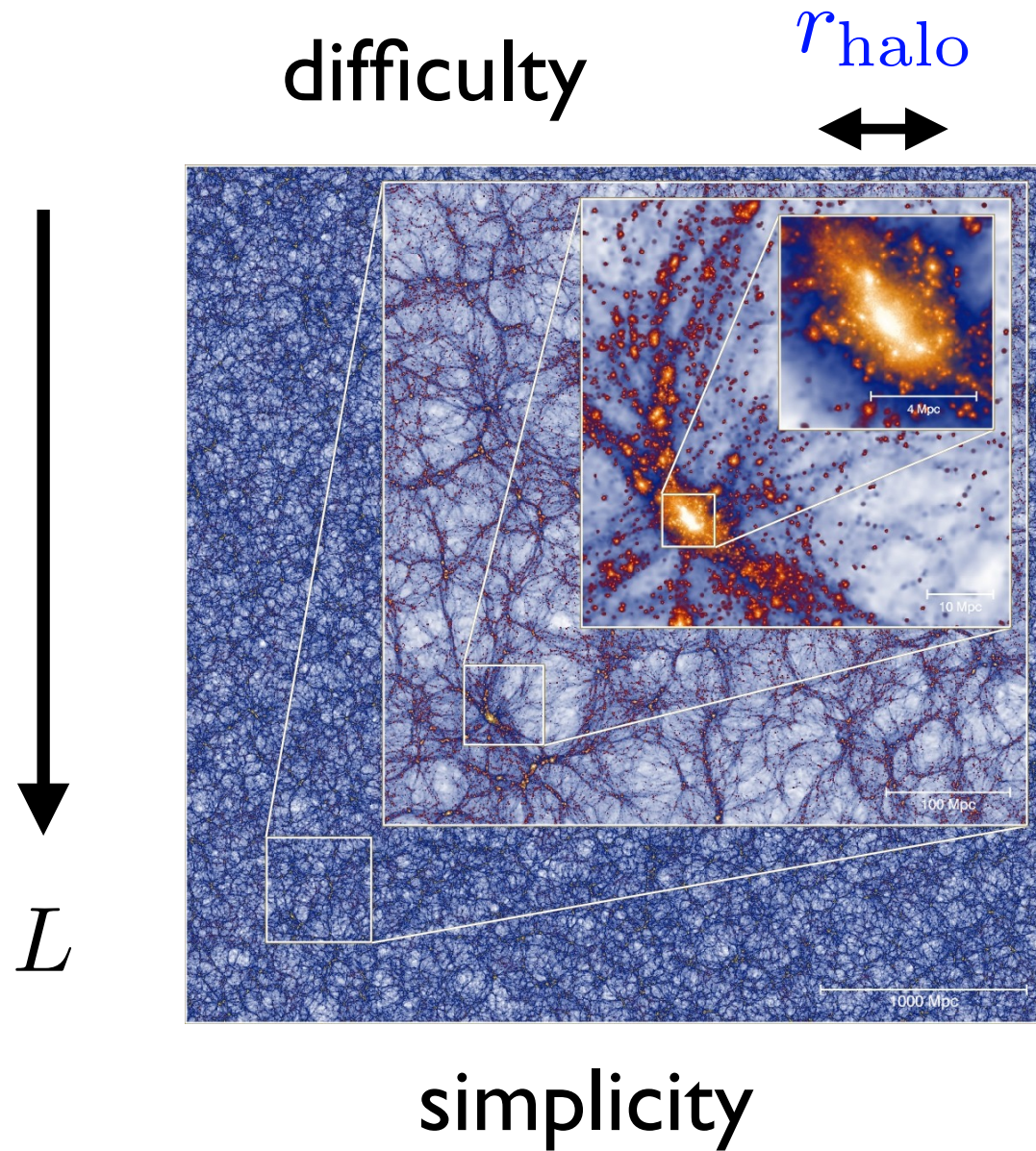
$$\delta m_H^2 = \text{---} \bigcirc \text{---} + \text{---} \text{ (irregular blob) } \text{---} + \text{---} \bigcirc \text{---}$$

© Donoghue, Golowich, Holstein SM textbook

● Cosmology: ensemble of virtual Universes

$$\delta P(k) = \text{---} \bigcirc \text{---} + \text{---} \bigcirc \text{---}$$

# Large scale structure perturbation theory



● Coarse-grained fields:

$$\frac{\partial u^i}{\partial t} + aH \partial u^i + u^j \nabla_j u^i = -\nabla^i \Phi - \frac{1}{\rho} \partial^j \sigma_{ij}$$

● Dim. analysis + symmetries:

$$\sigma_{ij} = c_s^2 \rho \delta_{ij} + \dots$$

● Caveats: time evolution is slow  $\sim$  age of Universe

● Radiation effects can make  $R_{\text{gal}}$  very large

Mirbabayi, Schmidt ++ (2014), ++

Cabass, Schmidt (2018)