## The S-Matrix and boundary correlators in flat space

Amplítudes 2024<br>Diksha Jain<br>TIFR Mumbaí<br>2311.03443

In Collaboration with Suman Kundu, Shiraz Minwalla, Onkar Parrikar, Siddharth Prabhu, Pushkal Shrivastava.

## Set up

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We ask what information does the path integral carry?

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Difference: In AdS $\phi_{0}$ couples to a boundary operator of dimension $\Delta$ and the corresponding path integral has independent description in dual theory.

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- However, most of our results can be generalised to arbitrary boundary surface (We consider null boundary as another example and make some comments about relation to C(FT).

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- S-matrix unitarity provides non-trivial constraint on this path integral.
- We conjecture that the flat space wave functional and the S-matrix are related by analytic continuation.
- We analysed the analytic structure of $G_{\text {brry }}$ in position space both for massive and massless particles.
- For massless particles, $G_{\text {bdry }}$ exhibits features like bulk point síngularity (and it's generalisations) whose coefficient encode the flat space S-matrix.

S-matrix as boundary observable

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S\left(\left\{p_{i}\right\},\left\{q_{j}\right\}\right)=\prod_{i=1}^{n} \int_{M} d^{d+1} x_{i} \sqrt{g\left(x_{i}\right)} f_{p_{i}}\left(x_{i}\right)\left(\nabla_{i}^{2}-m^{2}\right) \prod_{j=1}^{m} \int_{M} d^{d+1} y_{j} \sqrt{g\left(y_{j}\right)} \bar{f}_{q_{j}}\left(y_{j}\right)\left(\nabla_{j}^{2}-m^{2}\right) G\left(\left\{x_{i} y_{i}\right\}\right)
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\left.\int_{M} d^{d+1} x \sqrt{g(x)} f_{p}(x)\left(\nabla^{2}-m^{2}\right) \phi(x)=\int_{M} d^{d+1} x \partial_{\mu}\left(f_{p}(x) \sqrt{g(x)} \partial_{\mu} \phi(x)-\phi(x) \sqrt{g(x)} \partial_{\mu} f_{p}(x)\right)\right)
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=\prod_{i=1}^{n} \int_{B} d^{d} x_{i} \sqrt{h} n^{\mu_{i}}\left(f_{p_{i}} \partial_{\mu_{i}}-\partial_{\mu_{i}} f_{p_{i}}\right) \prod_{j=1}^{m} \int_{B} d^{d} y_{j} \sqrt{h} n^{\mu_{j}}\left(\bar{f}_{q_{j}} \partial_{\mu_{j}}-\partial_{\mu_{j}} \bar{f}_{q_{j}}\right) G\left(\left\{x_{i}, y_{j}\right\}\right)
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Z\left[\beta_{0}\right] & =\int_{\left.\phi\right|_{B}=\beta_{0}}[D \phi] e^{-S[\phi]} \\
G_{\text {bdry }}\left(\left\{x_{i}\right\},\left\{y_{i}\right\}\right) & =\left.\prod_{i=1}^{n} \frac{\delta}{\delta \beta_{0}\left(x_{i}\right)} \prod_{j=1}^{m} \frac{\delta}{\delta \beta_{0}\left(y_{j}\right)} Z\left[\beta_{0}\right]\right|_{\beta_{0}=0}
\end{aligned}
$$

Using Hamilton-Jacobí:

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G_{\text {bdry }}=\int_{\left.\phi\right|_{B}=0}[D \phi] \prod_{i=1}^{n} n^{\mu_{i}} \partial_{\mu_{i}} \phi\left(x_{i}\right) \prod_{j=1}^{m} n^{\mu_{j}} \partial_{\mu_{j}} \phi\left(y_{j}\right) e^{-S[\phi]}
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\left(f_{p_{i}} n^{\mu_{i}} \partial_{\mu_{i}} \phi\left(x_{i}\right)-\phi\left(x_{i}\right) n^{\mu_{i}} \partial_{\mu_{i}} f_{p_{i}}\right)
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$$
S_{E}\left(\left\{p_{i}\right\},\left\{q_{j}\right\}\right)=\int \prod_{i=1}^{n} d^{d} x_{i} f_{p_{i}}\left(x_{i}\right) \prod_{j=1}^{m} d^{d} y_{j} \bar{f}_{q_{j}}\left(y_{j}\right) G_{\text {bdry }}\left(\left\{x_{i}\right\},\left\{y_{j}\right\}\right)
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- Take the large T limit first and then analytically continue to Lorentizian space.


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- We conclude that the Dirichlet path integral serves as a generating functional for S-matrices.


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- AFS in 1974 made a similar statement.
[Arefeva, Faddeev \& Slavnov '74]
- But they considered Lorentzian path integral as a functional of "positive energy data in the past and negative energy data in future"- "In-Out Problem".


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- Such diagrams are different in Dirichlet and In-Out problem. In the Dirichlet case, we obtain vacuum wavefunction in $x$-basis and in the latter case, in coherent state basis.

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\psi\left(\left\{\bar{\beta}_{k}\right\}\right) \approx\left(-\int \prod_{i} d^{d} k_{i} \frac{i \lambda}{\sum_{i} \omega_{i}} \delta^{d}\left(\sum_{i} \vec{k}_{i}\right) \prod_{i} \bar{\beta}_{k_{i}}\right)
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$S^{\dagger} S=\rrbracket$

$$
\bar{\beta}=T^{\top} \quad \bar{\beta}^{\prime}
$$


$\int \mathscr{D} \bar{\beta}^{\prime} \mathscr{D} \bar{\beta} \exp \left(-\int \frac{d^{d} k}{(2 \pi)^{d}} 2 \omega_{k} \bar{\beta}_{-\bar{k}}^{*} \bar{\beta}_{\bar{k}}^{\prime}\right) Z^{*}[\beta, \bar{\beta}] \mathrm{C}\left[\beta^{\prime}, \bar{\beta}\right]=\exp \left(\int \frac{d^{d} p}{(2 \pi)^{d}} 2 \omega_{\bar{p}} \bar{\beta}_{\bar{p}}^{*} \beta \beta_{-\bar{p}}\right)$

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[J. Maldacena, D. Simmons-Duffin and A. Zhiboedov '17]

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## Singularity

The equation for pinch-off for $G_{\text {bdry }}$ can be phrased in terms of the distance matrix.

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Singularity appears when $N_{i j}$ has a zero eigenvalue with a positive eigenvector.

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\begin{array}{cll}
c=1 & \text { if } & \mathrm{m} \leq \mathrm{D}+1 \\
c=m-D & \text { if } & \mathrm{m}>\mathrm{D}+1
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Same as co-dimension of singularity.

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- The location of these singularities can be characterised in terms of zero eigenvalues of the boundary distance matrix: $N_{i j}=\left(x_{i}-x_{j}\right)^{2}$.
- The residue at these singularities contain the information about flat space $S$-matrix.

Two ways to extract S-matrix from $G_{\text {bdry }}$ :

1. Multiply with mode functions and integrate (essentially Fourier transform).
2. The coefficient of singularity of $G_{\text {bdry }}$ is the $S$ - matrix.

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[S. Banerjee '24]


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G(x, y)=C \frac{e^{-i m(x-y)^{2}}}{\left((x-y)^{2}\right)^{\frac{D-1}{4}}}
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$$
G_{\mathrm{bdry}} \approx \prod_{i=1}^{n}\left(\left(\frac{m_{i}}{2 \pi}\right)^{\frac{D-1}{2}} \frac{(-T-t)}{d_{i}^{i \frac{D+1}{2}}} e^{-i m_{i} d_{i}^{\text {in }}}\right) \prod_{i=1}^{m}(\text { out }) S\left(\frac{m_{i}\left(\vec{x}_{i}^{\text {in }}-\vec{y}\right)}{d_{i}^{\text {in }}}, \frac{m_{i}\left(\vec{x}_{i}^{\text {out }}-\vec{y}\right)}{d_{i}^{\text {out }}}\right)
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Holographic Renormalization is non-local!!

## Results/Outline

- S-matrix can be thought of as a boundary observable and can be computed using "Path integral as a functional of boundary values".
- S-matrix unitarity provides non-trivial constraint on this path integral.
- We argue that the flat space wave functional and the $S$-matrix are related by analytic continuation.
- We also analyse properties of $G_{\text {brry }}$ in position space both for massive and massless particles.
- For massless particles, $G_{\text {bdry }}$ exhibits features like bulk poínt singularity whose coefficient encode the flat space $S$-matrix.


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## ధస్లృలదదగగతు

நன்ली ड्रणाइा पंतहाव येదपषाथ நனறி றヘ円 ধনাবাদ Thank you फた గ్వantaxem धन्यवाद

