The S-Matrix and boundary correlators in flat space

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Set up







• We compute the path integral as a function of these boundary values.



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We ask what information does the path integral carry?







Ads/CFT



Ads/CFT $\phi(z,x) \sim z^{d-\Delta} \phi_0(x) + z^{\Delta} \phi_1(x)$



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 $\phi(z,x) \sim z^{d-\Delta}\phi_0(x) + z^{\Delta}\phi_1(x)$ Fix the value of growing mode (ϕ_0) at boundary $z = \epsilon$. Path integral as a function of (ϕ_0) carries information about the bulk dynamics.



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Difference: In AdS ϕ_0 couples to a boundary operator of dimension Δ and the corresponding path integral has independent description in dual theory.

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• However, most of our results can be generalised to arbitrary boundary surface (We consider null boundary as another example and make some comments about relation to CCFT).







• We provide a precise relationship between the flat space S-matrix and the "Path integral as a functional of boundary values".

Results/Outline



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• For massless particles, G_{bdry} exhibits features like bulk point singularity (and it's generalisations) whose coefficient encode the flat space S-matrix.





 $S(\{p_i\},\{q_j\}) = \prod_{i=1}^n \int_M d^{d+1}x_i \sqrt{g(x_i)} f_{p_i}(x_i) (\nabla_i^2 - m^2) \prod_{j=1}^m \int_M d^{d+1}y_j \sqrt{g(y_j)} \bar{f}_{q_j}(y_j) (\nabla_j^2 - m^2) G(\{x_i, y_i\})$



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 $\int_{M} d^{d+1}x \sqrt{g(x)} f_p(x) (\nabla^2 - m^2) \phi(x) = \int_{M} d^{d+1}x \,\partial_\mu \left(f_p(x) \sqrt{g(x)} \partial_\mu \phi(x) - \phi(x) \sqrt{g(x)} \partial_\mu f_p(x)) \right)$



 $= \prod_{i=1}^{n} \int_{B} d^{d}x_{i} \sqrt{h} n^{\mu_{i}} (f_{p_{i}}\partial_{\mu_{i}} - \partial_{\mu_{i}}f_{p_{i}}) \prod_{i=1}^{m} \int_{B} d^{d}y_{j} \sqrt{h} n^{\mu_{j}} (\bar{f}_{q_{j}}\partial_{\mu_{j}} - \partial_{\mu_{j}}\bar{f}_{q_{j}}) G(\{x_{i}, y_{j}\})$

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We can relate S-matrix (Euclidean version) to the bulk Euclidean pathintegral with specified boundary conditions on a boundary surface B.



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 $G_{\text{bdry}}(\{x_i\}, \{y_i\}) = \prod_{i=1}^{n} \frac{\delta}{\delta\beta_0(x_i)} \prod_{i=1}^{m} \frac{\delta}{\delta\beta_0(y_i)} Z[\beta_0] \Big|_{\beta_0 = 0}$

Path Integral

We can relate S-matrix (Euclidean version) to the bulk Euclidean pathintegral with specified boundary conditions on a boundary surface B.

 $Z[\beta_0] = \int_{\phi|_B = \beta_0} [D\phi] e^{-S[\phi]}$




$$G_{\rm bdry} = \int_{\phi|_B=0} [D\phi] \int_{i=1}^{n} D\phi_{\rm bdry} \int_{i=1}^{n} D\phi_{\rm bdry}$$

 $\prod_{i=1}^{n} n^{\mu_i} \partial_{\mu_i} \phi(x_i) \prod_{j=1}^{m} n^{\mu_j} \partial_{\mu_j} \phi(y_j) e^{-S[\phi]}$



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 $(f_{p_i}n^{\mu_i}\partial_{\mu_i}\phi(x_i) - \phi(x_i)n^{\mu_i}\partial_{\mu_i}f_{p_i})$



$$G_{\text{bdry}} = \int_{\phi|_B=0} [D\phi] \begin{bmatrix} D\phi \end{bmatrix} \prod_{i=1}^{n} f_{i}$$

$$S_E(\{p_i\}, \{q_j\}) = \int_{i=1}^n d^{d}$$

 $\prod_{j=1}^{m} n^{\mu_i} \partial_{\mu_i} \phi(x_i) \prod_{j=1}^{m} n^{\mu_j} \partial_{\mu_j} \phi(y_j) e^{-S[\phi]}$

 ${}^{d}x_{i}f_{p_{i}}(x_{i})\prod_{j=1}^{m}d^{d}y_{j}\bar{f}_{q_{j}}(y_{j})G_{bdry}(\{x_{i}\},\{y_{j}\})$





Euclidean Dirichlet problem

Lorentzian Dirichlet problem but with a twist



Euclidean Dirichlet problem

space.

Lorentzian Dirichlet problem but with a twist • Take the large T limit first and then analytically continue to Lorentizian



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$$S(\{p_i\}, \{q_j\}) = \int_{i=1}^{n} d^d x$$

Lorentzian Dirichlet problem but with a twist • Take the large T limit first and then analytically continue to Lorentizian

 $x_i f_{p_i}(x_i) \prod d^d y_j \bar{f}_{q_j}(y_j) G^L_{bdry}(\{x_i\}, \{y_j\})$ j=1



Euclidean Dirichlet problem

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$$S(\{p_i\}, \{q_j\}) = \int_{i=1}^{n} d^d x$$

functional for S-matrices.

Lorentzian Dirichlet problem but with a twist • Take the large T limit first and then analytically continue to Lorentizian

 $x_i f_{p_i}(x_i) \prod d^d y_j \bar{f}_{q_i}(y_j) G^L_{bdry}(\{x_i\}, \{y_j\})$ j=1• We conclude that the Dirichlet path integral serves as a generating







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[Arefeva, Faddeev & Slavnov '74]



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• AFS in 1974 made a similar statement.

 Faddeev & Slavnov '74]
 But they considered Lorentzian path integral as a functional of "positive energy data in the past and negative energy data in future" — "In-Out Problem".





• The quantities which receive contributions from deep inside the bulk give the same answer in both Dirichlet and In-Out problem.



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• The Path integral $Z[\beta]$ carries much more information than just the S-matrix.



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 Such diagrams are different in Dirichlet and In-Out problem. In the Dirichlet case, we obtain vacuum wavefunction in x-basis and in the latter case, in coherent state basis.







 $\psi(\{\bar{\beta}_k\}) \approx \left(-\int \prod_i d^d k_i \, \frac{i\lambda}{\sum_i \omega_i} \delta^d \left(\sum_i \bar{k}_i\right) \prod_i \bar{\beta}_{k_i}\right)$



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Wavefunction \implies S-matrix

• We found that at tree level, the wave function contains a pole in $\sum \omega_i$.

[P. Beníncasa '18,...]





Unitarity





Unitarity $S^{\dagger}S = 0$



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Coherent State Interpretation



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 $G_{\text{bdry}} = \lambda \int d^{d+1}y \prod_{i=1}^{n} G_{\partial B}(x_i, y)$





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Near Singularity

G_{bdry}: massless particles





$$G_{\text{bdry}} = \lambda \int d^{d+1}y \prod_{i=1}^{n} G_{i=1}$$

Near Singularity

 $G_{\partial B}(x_i, y) = (2n \cdot \nabla G(x, y)) \Big|_{x \to x_i}$

sless particles

 $f_{\partial B}(x_i, y)$









G_{bdry}: massless particles $G(x, y) = \int d^{d+1}y \prod_{i=1}^{n} \left(\frac{1}{\left((x_i - y)^2 - i\epsilon \right)^{\frac{D-2}{2}}} \right)$



 $G(x, y) = \int d^{d+1}y \prod_{i=1}^{n} d^{i-1}y \prod_{i=1}$

G(x, y) has pole-type singularities whenever $(x_i - y)^2 = 0$.

$$\int_{1}^{1} \left(\frac{1}{(x_i - y)^2 - i\epsilon} \right)^{\frac{D-2}{2}}$$



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[J.

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Maldacena, D. Simmons-Duffin and A. Zhiboedov





The equation for pinch-off for G_{bdry} can be phrased in terms of the distance matrix.



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Singularity

The equation for pinch-off for G_{bdry} can be phrased in terms of the

 $N_{ij} = (x_i - x_j)^2 = \left((x_i - y) - (x_j - y) \right)^2$



The equation for pinch-off for $G_{\rm bdry}$ can be phrased in terms of the distance matrix.

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t y s.t. $(x_i - y)^2 = 0$



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 $N_{ij} = (x_i - x_j)^2 = (0)$ Assuming there exists a bulk point $N_{ij} = 2(x_i - x_j)^2 = (1)$

$$(x_{i} - y) - (x_{j} - y) \Big)^{2}$$

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Pinch-off/Momentum conservation:



$$(x_{i} - y) - (x_{j} - y) \Big)^{2}$$

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Pinch-off/Momentum conservation:



Singularity

$$(x_i - y) - (x_j - y) \Big)^2$$

tys.t. $(x_i - y)^2 = 0$

 $\sum \omega_i N_{ij} = 0$

Singularity appears when N_{ii} has a zero eigenvalue with a positive eigenvector.



Co-dimension of Singularity



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c = 1 if $m \le D + 1$ $c = m - D \quad \text{if} \quad m > D + 1$





insertions, does G_{bdry} receive contributions from one S-matrix or many?

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Again depends on number of insertions and dimension of spacetime.

When m ≤ D + 1, only one S-matrix.
When m > D + 1, G_{bdry} receives contributions from m − D, S-matrices.



Q2: Given a set of boundary points $\{x_i\}$ such that G_{bdry} is singular for those insertions, does G_{bdry} receive contributions from one S-matrix or many?

Again depends on number of insertions and dimension of spacetime.

When m ≤ D + 1, only one S-matrix.
When m > D + 1, G_{bdry} receives contributions from m − D, S-matrices.

Same as co-dimension of singularity.







G_{bdry}: massless particles • We find that for massless particles, $G_{bdry}(x_i)$ (at tree level) is an analytic function in the space of boundary insertions with pole type singularities.



 $(c \ge 1)$ in the space of boundary insertions.

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 The residue at these singularities contain the information about flat

space S-matrix.




Two ways to extract S-matrix from G_{bdry} : 1. Multiply with mode functions and integrate (essentially Fourier transform). 2. The coefficient of singularity of G_{bdry} is the S- matrix.





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[S. Banerjee '24]





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 $G(x, y) = C - \frac{e^{-im(x-y)^2}}{D}$ $((x-y)^2)^{\frac{D-1}{4}}$

G_{bdry}: massive particles We computed boundary correlates in position space at tree level for massive









This equation gives momentum conservation at the bulk point y.

 $\sum_{i} m_i \frac{(x_i - y)^{\mu}}{d_i(x_i, y)} = 0$





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$$G_{\text{bdry}} \approx \prod_{i=1}^{n} \left(\left(\frac{m_i}{2\pi} \right)^{\frac{D-1}{2}} \frac{(-T-t)}{d_i^{in^{\frac{D+1}{2}}}} e^{-t} \right)^{\frac{D-1}{2}}$$

 $-im_{i}d_{i}^{in}\prod_{i=1}^{m}(\text{out})S\left(\frac{m_{i}(\vec{x}_{i}^{in}-\vec{y})}{d_{i}^{in}},\frac{m_{i}(\vec{x}_{i}^{out}-\vec{y})}{d_{i}^{out}}\right)$





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Holographic Renormalization is non-local!!



Results/Outline

- related by analytic continuation.
- and massless particles.
- For massless particles, $G_{\rm bdry}$ exhibits features like bulk point singularity whose coefficient encode the flat space S-matrix.

• S-matrix can be thought of as a boundary observable and can be computed using "Path integral as a functional of boundary values". • S-matrix unitarity provides non-trivial constraint on this path integral. • We argue that the flat space wave functional and the S-matrix are

• We also analyse properties of $G_{\rm bdry}$ in position space both for massive





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ಧನ್ಯವಾದಗಳು आशार துரு येतरान् भिन्नज्व हार्जा पेतरान् भिन्नज्व भग्रावाम Thank you ຜູ້ສຸສາຍສາຍ ຊິຽຽ